

A space-time multiscale mortar mixed finite element method for parabolic equations

Manu Jayadharan, Michel Kern, **Martin Vohralík**, and Ivan Yotov

Inria Paris & Ecole des Ponts

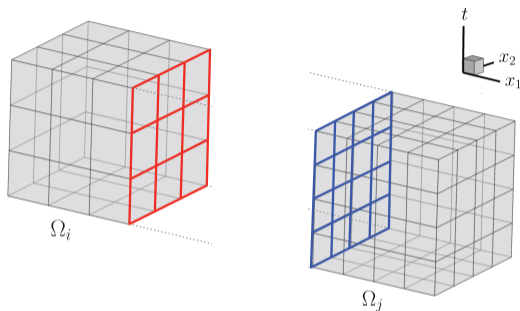
Oberwolfach, February 7, 2022

The Inria logo is written in a red, cursive script.

Outline

- 1 Introduction
- 2 Space-time multiscale mortar mixed finite element method
 - Continuous setting
 - Discrete setting
 - Space-time multiscale mortar mixed finite element method
 - Existence, uniqueness, and stability
- 3 Reduction to an interface problem
- 4 Numerical experiments
- 5 Conclusions and future directions

A space-time multiscale mortar mixed finite element method



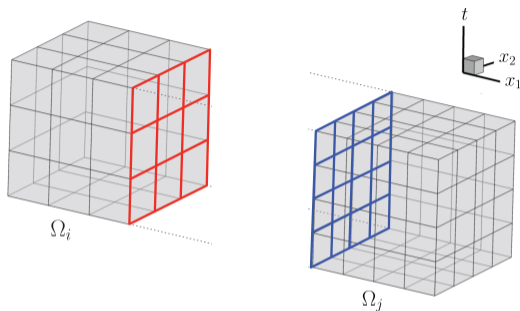
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Concepts

- computational domain Ω (polytope)
- partition of Ω into non-overlapping polytopal **subdomains** $\bar{\Omega} = \cup \bar{\Omega}_j$

different space and time discretizations on subdomains (local time stepping),
coupling through **multiscale space-time mortars, space-time DD**

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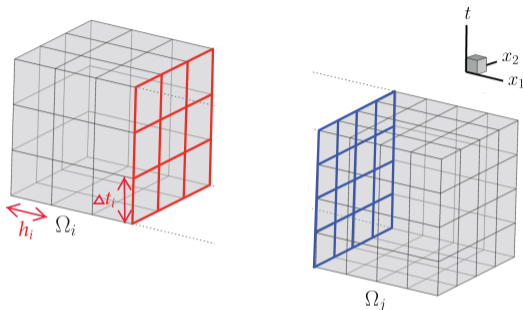
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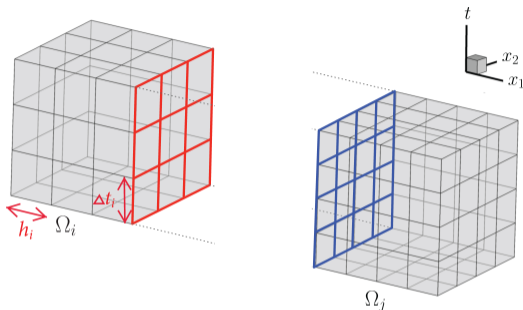
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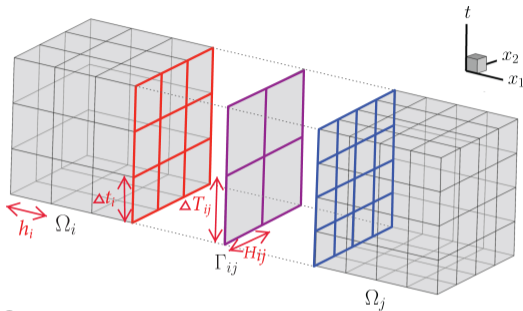
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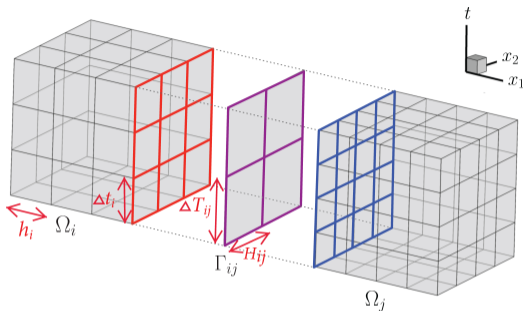
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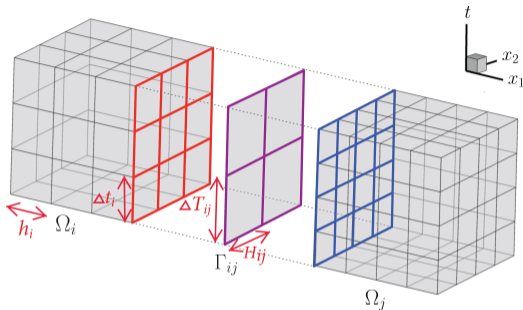
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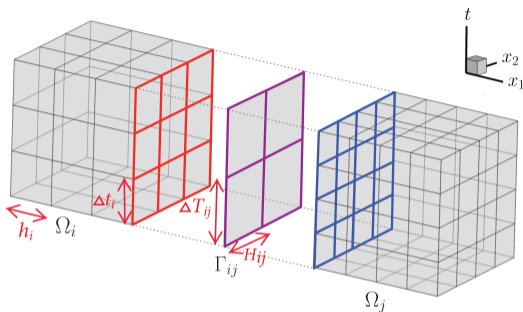
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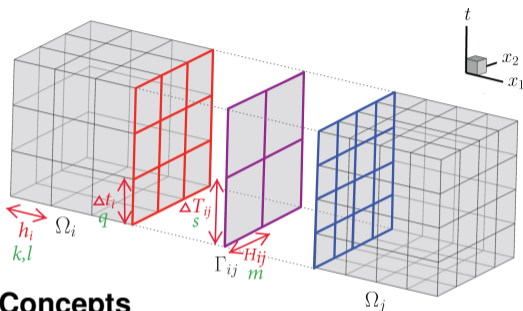
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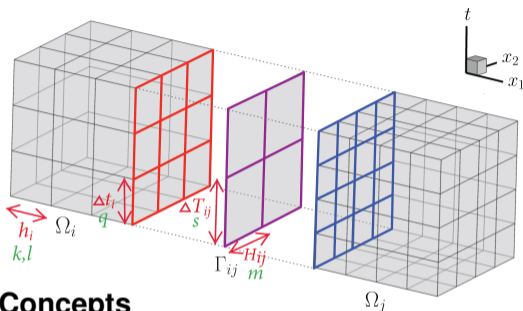
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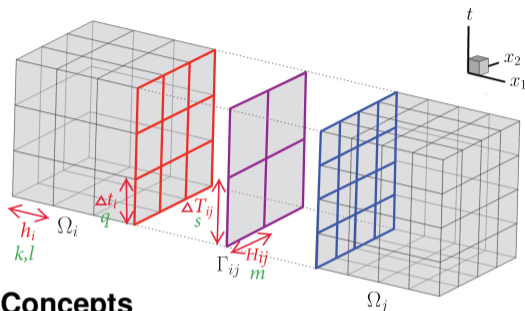
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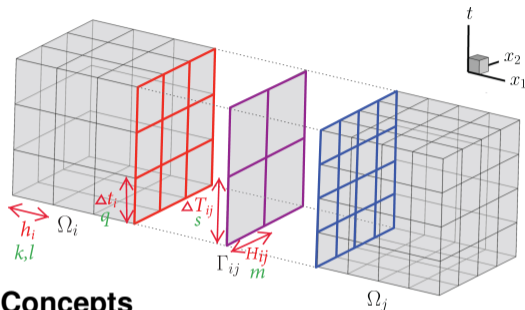
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Context – steady case

Hybridized formulation of mixed finite element methods

- Fraeijs de Veubeke (1960's)
- Arnold and Brezzi (1985)
- Arbogast and Chen (1995)
 - interface mesh given by the neighboring subdomains
 - interface: same mesh and same polynomial degree as in the subdomains
 - hybridized and initial problems equivalent

Multiscale mortar mixed finite element method

- Arbogast, Pencheva, Wheeler, and Yotov (2007)
 - independent interface mesh
 - typically coarser but one employs polynomials of higher degree
 - (multiscale) weak continuity of the normal flux component over the interfaces between subdomains
 - efficient parallelization via a non-overlapping domain decomposition algorithm reducing to an interface problem
- ... related to numerous other multiscale approaches

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Context – unsteady case

Local time stepping for parabolic problems

- Ewing, Lazarov, and Vassilevski (1990), Delpopolo Carciopolo, Cusini, Formaggia, and Hajibeygi (2020), ...

Domain decomposition methods with local time stepping

- Dawson, Du, and Dupont (1991), Yu (2001), Gaiffe, Glowinski, and Masson (2002), Faucher and Combescure (2003), Gander and Halpern (2007), Nakshatrala, Nakshatrala, and Tortorelli (2009), Hager, Hauret, Le Tallec, and Wohlmuth (2012), Kheriji, Masson, and A. Moncorgé (2015), Krause and Krause (2016), Arshad, Park, and Shin (2021), ...

Space-time domain decomposition

- Halpern, Japhet, and Szeftel (2012), Hoang, Jaffré, Japhet, Kern, and Roberts (2013), Gander, Kwok, and Mandal (2016), ...

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- Lions, Maday, and Turinici (2001), Gander and Vandewalle (2007), Falgout, Friedhoff, Kolev, MacLachlan, and Schroder (2014), Gander and Neumüller (2016), ...

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Setting

The heat equation

Find $p : \Omega \times [0, T] \rightarrow \mathbb{R}$ and $\mathbf{u} : \Omega \times [0, T] \rightarrow \mathbb{R}^d$ such that

$$\begin{aligned} \frac{\partial p}{\partial t} + \nabla \cdot \mathbf{u} &= q && \text{in } \Omega \times (0, T], \\ \mathbf{u} &= -K \nabla p && \text{in } \Omega \times (0, T], \\ p &= 0 && \text{on } \partial\Omega \times (0, T], \\ p &= p_0(x) && \text{on } \Omega. \end{aligned}$$

- $q \in L^2(0, T; L^2(\Omega))$, $p_0 \in H_0^1(\Omega)$, $\nabla \cdot K \nabla p_0 \in L^2(\Omega)$
- K : time-independent, uniformly bounded, symmetric, and positive definite

Weak solution

Find $\mathbf{u} \in L^2(0, T; \mathbf{H}(\text{div}; \Omega))$, $p \in H^1(0, T; L^2(\Omega))$ s.t. $p(\cdot, 0) = p_0$ & a.e. in $(0, T)$,

$$\begin{aligned} (K^{-1} \mathbf{u}, \mathbf{v})_{\Omega} - (p, \nabla \cdot \mathbf{v})_{\Omega} &= 0 \quad \forall \mathbf{v} \in \mathbf{H}(\text{div}; \Omega), \\ (\partial_t p, w)_{\Omega} + (\nabla \cdot \mathbf{u}, w)_{\Omega} &= (q, w)_{\Omega} \quad \forall w \in L^2(\Omega). \end{aligned}$$

- actually also $\mathbf{u} \in L^\infty(0, T; \mathbf{L}^2(\Omega))$ and $p \in H^1(0, T; H_0^1(\Omega))$

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Space-time subdomains, forms, and spaces

Space-time subdomains

$$\Omega^T := \Omega \times (0, T), \Omega_i^T := \Omega_i \times (0, T), \Gamma_i^T := \Gamma_i \times (0, T), \Gamma_{ij}^T := \Gamma_{ij} \times (0, T)$$

Space-time bilinear forms

$$a^T(\mathbf{u}, \mathbf{v}) := \int_0^T \sum_i (K^{-1} \mathbf{u}, \mathbf{v})_{\Omega_i}, \quad b^T(\mathbf{v}, w) := - \int_0^T \sum_i (\nabla \cdot \mathbf{v}, w)_{\Omega_i},$$

$$b_\Gamma^T(\mathbf{v}, \mu) := \int_0^T \sum_i \langle \mathbf{v} \cdot \mathbf{n}_i, \mu \rangle_{\Gamma_i}$$

Tensor product space-time spaces on each space-time subdomain Ω_i^T

$$\mathbf{V}_{h,i}^{\Delta t} := \underbrace{\mathbf{V}_{h,i}}_{\text{MFE spaces}} \otimes \underbrace{W_i^{\Delta t}}_{\text{discontinuous pw polynomials in time}}, \quad W_{h,i}^{\Delta t} := \underbrace{W_{h,i}}_{\text{discontinuous pw polynomials in space}} \otimes W_i^{\Delta t}$$

$$\Lambda_{H,ij}^{\Delta T} := \underbrace{\Lambda_{H,ij} \otimes \Lambda_{ij}^{\Delta T}}_{\text{continuous or discontinuous pw polynomials in space and in time}}$$

Space-time subdomains, forms, and spaces

Space-time subdomains

$$\Omega^T := \Omega \times (0, T), \Omega_i^T := \Omega_i \times (0, T), \Gamma_i^T := \Gamma_i \times (0, T), \Gamma_{ij}^T := \Gamma_{ij} \times (0, T)$$

Space-time bilinear forms

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Global spaces and time stepping

Global space-time finite element spaces

$$\mathbf{V}_h^{\Delta t} := \bigoplus \mathbf{V}_{h,i}^{\Delta t}, \quad W_h^{\Delta t} := \bigoplus W_{h,i}^{\Delta t}, \quad \Lambda_H^{\Delta T} := \bigoplus \Lambda_{H,ij}^{\Delta T}$$

Space of velocities with space-time weakly continuous normal components

$$\mathbf{V}_{h,0}^{\Delta t} = \left\{ \mathbf{v} \in \mathbf{V}_h^{\Delta t} : b_{\Gamma}^T(\mathbf{v}, \mu) = 0 \quad \forall \mu \in \Lambda_H^{\Delta T} \right\}$$

Discontinuous Galerkin time stepping

$$\int_0^T \tilde{\partial}_t \varphi \phi = \sum_{k=1}^{N_i} \int_{t_j^{k-1}}^{t_j^k} \partial_t \varphi \phi + \sum_{k=1}^{N_i} [\varphi]_{k-1} \phi_{k-1}^+,$$

with

$$[\varphi]_k = \varphi_k^+ - \varphi_k^-, \quad \varphi_k^+ = \lim_{t \rightarrow t_j^{k,+}} \varphi, \quad \varphi_k^- = \lim_{t \rightarrow t_j^{k,-}} \varphi$$

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- 2 Space-time multiscale mortar mixed finite element method
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Space-time multiscale mortar mixed finite element method

Definition (Space-time multiscale mortar mixed finite element method)

Find $\mathbf{u}_h^{\Delta t} \in \mathbf{V}_h^{\Delta t}$, $p_h^{\Delta t} \in W_h^{\Delta t}$, and $\lambda_H^{\Delta T} \in \Lambda_H^{\Delta T}$ such that

$$\begin{aligned} a^T(\mathbf{u}_h^{\Delta t}, \mathbf{v}) + b^T(\mathbf{v}, p_h^{\Delta t}) + b_\Gamma^T(\mathbf{v}, \lambda_H^{\Delta T}) &= 0 & \forall \mathbf{v} \in \mathbf{V}_h^{\Delta t}, \\ (\tilde{\partial}_t p_h^{\Delta t}, w)_{\Omega^T} - b^T(\mathbf{u}_h^{\Delta t}, w) &= (q, w)_{\Omega^T} & \forall w \in W_h^{\Delta t}, \\ b_\Gamma^T(\mathbf{u}_h^{\Delta t}, \mu) &= 0 & \forall \mu \in \Lambda_H^{\Delta T}. \end{aligned}$$

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Assumptions

Assumption (Mortar grids)

For C independent of the spatial mesh sizes h and H as well as of the temporal mesh sizes Δt and ΔT , there holds

$$\forall \mu \in \Lambda_H, \forall i, j, \quad \|\mu\|_{\Gamma_{ij}} \leq C(\|\mathcal{Q}_{h,i} \mu\|_{\Gamma_{ij}} + \|\mathcal{Q}_{h,j} \mu\|_{\Gamma_{ij}}),$$

$$\forall i, j, \quad \Lambda_{ij}^{\Delta T} \subset W_i^{\Delta t} \cap W_j^{\Delta t}.$$

Comments

- $\Lambda_H := \bigoplus \Lambda_{H,ij}$
- $\mathcal{Q}_{h,i} : L^2(\partial\Omega_i) \rightarrow \mathbf{V}_{h,i} \cdot \mathbf{n}_i$ is the L^2 -orthogonal projection
- the spatial mortar assumption as previously: in particular satisfied with $C = \frac{1}{2}$ when $\mathcal{T}_{H,ij}$ is a coarsening of both $\mathcal{T}_{h,i}$ and $\mathcal{T}_{h,j}$ on Γ_{ij} & the space $\Lambda_{H,ij}$ consists of discontinuous pw polynomials contained in $\mathbf{V}_{h,i} \cdot \mathbf{n}_i$ and $\mathbf{V}_{h,j} \cdot \mathbf{n}_j$ on Γ_{ij} ; in general, it requests the mortar space Λ_H to be sufficiently coarse
- the temporal mortar assumption: control of the mortar by the subdomain time discretizations; mortar time discretization is a coarsening of each subdomain

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Two inf-sup inequalities

Lemma (Discrete **divergence** inf-sup condition on $\mathbf{V}_{h,0}^{\Delta t}$)

Let the mortar assumptions hold. Then

$$\forall \mathbf{w} \in \mathbf{W}_h^{\Delta t}, \quad \sup_{0 \neq \mathbf{v} \in \mathbf{V}_{h,0}^{\Delta t}} \frac{\mathbf{b}^T(\mathbf{v}, \mathbf{w})}{\|\mathbf{v}\|_{L^2(0,T;\Pi_i \mathbf{H}(\text{div}; \Omega_i))}} \geq \beta \|\mathbf{w}\|_{L^2(0,T;L^2(\Omega))}.$$

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Existence, uniqueness, and stability

Theorem (Existence and uniqueness of the discrete solution, stability wrt data)

Let the mortar assumptions hold. Then the space-time multiscale mortar MFE method has a unique solution. Moreover,

$$\|p_h^{\Delta t}\|_{\text{DG}} + \|\mathbf{u}_h^{\Delta t}\|_{\Omega T} + \|p_h^{\Delta t}\|_{\Omega T} + \|\lambda_H^{\Delta T}\|_{\Gamma T} \leq C(\|q\|_{\Omega T} + \|\nabla \cdot K \nabla p_0\|_{\Omega}).$$

Comments

- $\|\varphi\|_{\text{DG}}^2 = \sum_i (\|\varphi_{N_i}^-\|_{\Omega_i}^2 + \sum_{k=1}^{N_i} \|[\varphi]_{k-1}\|_{\Omega_i}^2)$
- no control of divergence (shown later in a simplified setting)

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A priori error estimate

Theorem (A priori error estimate)

Let the mortar assumptions hold and let the weak solution be sufficiently smooth. Let the space and time meshes $\mathcal{T}_{h,i}$ and $\mathcal{T}_i^{\Delta t}$ be *quasi-uniform* and let $h \leq Ch_i$ and $\Delta t \leq C\Delta t_i$ for all i . Then

$$\begin{aligned} & \|p - p_h^{\Delta t}\|_{\text{DG}} + \|\mathbf{u} - \mathbf{u}_h^{\Delta t}\|_{\Omega_T} + \|p - p_h^{\Delta t}\|_{\Omega_T} + \|\lambda - \lambda_H^{\Delta T}\|_{\Gamma_T} \\ & \leq C \left(\sum_i \|\mathbf{u}\|_{H^{r_q}(0,T;H^{r_k}(\Omega_i))} (h^{r_k} + \Delta t^{r_q}) + \|\mathbf{u}\|_{H^{r_q}(0,T;H^{\tilde{r}_k+\frac{1}{2}}(\Omega))} (h^{\tilde{r}_k} H^{\frac{1}{2}} + \Delta t^{r_q}) \right. \\ & \quad \left. + \sum_i \|p\|_{W^{r_q,\infty}(0,T;H^{r_l}(\Omega_i))} \Delta t^{-\frac{1}{2}} (h^{r_l} + \Delta t^{r_q}) + \sum_{i,j} \|\lambda\|_{H^{r_s}(0,T;H^{r_m}(\Gamma_{ij}))} h^{-\frac{1}{2}} (H^{r_m} + \Delta T^{r_s}) \right), \\ & \underbrace{0 < r_k \text{ or } \tilde{r}_k \leq k + 1}_{\text{MFE space}}, \underbrace{0 \leq r_l \leq l + 1}_{\text{pw polys space}}, \underbrace{0 \leq r_q \leq q + 1}_{\text{pw polys time}}, \underbrace{0 \leq r_m \leq m + 1}_{\text{mortars space}}, \underbrace{0 \leq r_s \leq s + 1}_{\text{mortars time}}. \end{aligned}$$

A priori error estimate

Comments

- the term $h^{-\frac{1}{2}}(H^{r_m} + \Delta T^{r_s})$ appears from the discrete trace (inverse) inequality and is **suboptimal**; can be made comparable to the other error terms by choosing m and s **sufficiently large** (if the solution is sufficiently smooth)
- the term $\Delta t^{-\frac{1}{2}}(h^{r_l} + \Delta t^{r_q})$ is **suboptimal**
- **both improved** if a bound on $\|\nabla \cdot (\mathbf{u} - \mathbf{u}_h^{\Delta t})\|_{\Omega_j^T}$ is available, using the normal trace inequality for $\mathbf{H}(\text{div}; \Omega_j)$

Improved stability

Radau reconstruction operator

- for $\varphi(x, \cdot) \in W^{\Delta t}$, $\mathcal{I}\varphi(x, \cdot) \in H^1(0, T)$, $\mathcal{I}\varphi(x, \cdot)|_{(t^{k-1}, t^k)} \in P_{q+1}$, such that

$$\int_{t^{k-1}}^{t^k} \partial_t \mathcal{I}\varphi \phi = \int_{t^{k-1}}^{t^k} \partial_t \varphi \phi + [\varphi]_{k-1} \phi_{k-1}^+ \quad \forall \phi(x, \cdot) \in W^{\Delta t}$$

- thus equivalently, $\tilde{\partial}_t p_h^{\Delta t}$ replaced by $\partial_t \mathcal{I}p_h^{\Delta t}$:

$$(\partial_t \mathcal{I}p_h^{\Delta t}, w)_{\Omega T} - b^T(\mathbf{u}_h^{\Delta t}, w) = (q, w)_{\Omega T} \quad \forall w \in W_h^{\Delta t}$$

Theorem (Control of divergence)

Let $W_i^{\Delta t} = \Lambda_{ij}^{\Delta T} = W_j^{\Delta t}$ (same time discretization everywhere) hold. Then

$$\|\partial_t \mathcal{I}p_h^{\Delta t}\|_{\Omega T} + \|\nabla_h \cdot \mathbf{u}_h^{\Delta t}\|_{\Omega T} + \|\mathbf{u}_h^{\Delta t}\|_{\text{DG}} \leq C(\|q\|_{\Omega T} + \|\nabla \cdot K \nabla p_0\|_{\Omega}).$$

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Improved a priori error estimate

Theorem (Improved a priori error estimate)

Let the mortar space assumption and $W_i^{\Delta t} = \Lambda_{ij}^{\Delta T} = W_j^{\Delta t}$ hold and let the weak solution be sufficiently smooth. Then

$$\begin{aligned}
 & \| \mathbf{u}(t^N) - (\mathbf{u}_h^{\Delta t})_N^- \|_{\Omega} + \| \mathbf{u} - \mathbf{u}_h^{\Delta t} \|_{\Omega^T} + \| \nabla_h \cdot (\mathbf{u} - \mathbf{u}_h^{\Delta t}) \|_{\Omega^T} \\
 & \quad + \| p(t^N) - (p_h^{\Delta t})_N^- \|_{\Omega} + \| p - p_h^{\Delta t} \|_{\Omega^T} + \| \lambda - \lambda_H^{\Delta T} \|_{\Gamma^T} \\
 & \leq C \left(\sum_i \| \mathbf{u} \|_{W^{r_q, \infty}(0, T; \mathbf{H}^{r_k}(\Omega_i))} (h^{r_k} + \Delta t^{r_q}) + \| \mathbf{u} \|_{W^{r_q, \infty}(0, T; \mathbf{H}^{\tilde{r}_k + \frac{1}{2}}(\Omega))} (h^{\tilde{r}_k} H^{\frac{1}{2}} + \Delta t^{r_q}) \right. \\
 & \quad \left. + \sum_i \| p \|_{W^{r_q, \infty}(0, T; H^{r_l}(\Omega_i))} (h^{r_l} + \Delta t^{r_q}) + \sum_{i,j} \| \lambda \|_{H^{r_s}(0, T; H^{r_m}(\Gamma_{ij}))} (H^{r_m - \frac{1}{2}} + \Delta T^{r_s}) \right. \\
 & \quad \left. + \sum_i \| \mathbf{u} \|_{H^1(0, T; \mathbf{H}^{r_k}(\Omega_i))} h^{r_k} + \| \mathbf{u} \|_{H^1(0, T; \mathbf{H}^{\tilde{r}_k + \frac{1}{2}}(\Omega))} h^{\tilde{r}_k} H^{\frac{1}{2}} + \sum_{i,j} \| \lambda \|_{H^1(0, T; H^{r_m}(\Gamma_{ij}))} H^{r_m - \frac{1}{2}} \right), \\
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Decomposition of the solution

Decomposition of the solution

- $\mathbf{u}_h^{\Delta t} = \mathbf{u}_h^{\Delta t,*}(\lambda_H^{\Delta T}) + \bar{\mathbf{u}}_h^{\Delta t}$, $p_h^{\Delta t} = p_h^{\Delta t,*}(\lambda_H^{\Delta T}) + \bar{p}_h^{\Delta t}$
- for each Ω_i^T , $\bar{\mathbf{u}}_h^{\Delta t}|_{\Omega_i^T} \in \mathbf{V}_{h,i}^{\Delta t}$, $\bar{p}_h^{\Delta t}|_{\Omega_i^T} \in W_{h,i}^{\Delta t}$ is the solution to (zero Dirichlet data on the space-time interfaces and the prescribed source term q , initial data p_0 , and 0 boundary data on the external boundary)

$$a_i^T(\bar{\mathbf{u}}_h^{\Delta t}, \mathbf{v}) + b_i^T(\mathbf{v}, \bar{p}_h^{\Delta t}) = 0 \quad \forall \mathbf{v} \in \mathbf{V}_{h,i}^{\Delta t}$$

$$(\tilde{\partial}_t \bar{p}_h^{\Delta t}, w)_{\Omega_i^T} - b_i^T(\bar{\mathbf{u}}_h^{\Delta t}, w) = (q, w)_{\Omega_i^T} \quad \forall w \in W_{h,i}^{\Delta t}$$

- for a given $\mu \in \Lambda_H^{\Delta T}$, for each Ω_i^T , $\mathbf{u}_h^{\Delta t,*}(\mu)|_{\Omega_i^T} \in \mathbf{V}_{h,i}^{\Delta t}$, $p_h^{\Delta t,*}(\mu)|_{\Omega_i^T} \in W_{h,i}^{\Delta t}$ is the solution to (Dirichlet data μ , zero source term, initial data, and boundary data)

$$a_i^T(\mathbf{u}_h^{\Delta t,*}(\mu), \mathbf{v}) + b_i^T(\mathbf{v}, p_h^{\Delta t,*}(\mu)) = -\langle \mathbf{v} \cdot \mathbf{n}_i, \mu \rangle_{\Gamma_i^T} \quad \forall \mathbf{v} \in \mathbf{V}_{h,i}^{\Delta t},$$

$$(\tilde{\partial}_t p_h^{\Delta t,*}(\mu), w)_{\Omega_i^T} - b_i^T(\mathbf{u}_h^{\Delta t,*}(\mu), w) = 0 \quad \forall w \in W_{h,i}^{\Delta t}$$

- both above problems are posed in the individual space-time subdomains Ω_i^T and can thus be solved **in parallel** (no synchronization on time steps)

Space-time Steklov–Poincaré operator

Lemma (Equivalence)

The MMMFE method is **equivalent to**: find $\lambda_H^{\Delta T} \in \Lambda_H^{\Delta T}$ such that

$$-b_\Gamma^T(\mathbf{u}_h^{\Delta t,*}(\lambda_H^{\Delta T}), \mu) = b_\Gamma^T(\bar{\mathbf{u}}_h^{\Delta t}, \mu) \quad \forall \mu \in \Lambda_H^{\Delta T}.$$

Space-time Steklov–Poincaré operator

- $S : \Lambda_H^{\Delta T} \rightarrow \Lambda_H^{\Delta T}$

$$\langle S\lambda, \mu \rangle_{\Gamma T} := \sum_i \langle S_i \lambda, \mu \rangle_{\Gamma_i^T}, \quad \langle S_i \lambda, \mu \rangle_{\Gamma_i^T} := -\langle \mathbf{u}_h^{\Delta t,*}(\lambda) \cdot \mathbf{n}_i, \mu \rangle_{\Gamma_i^T} \quad \forall \lambda, \mu \in \Lambda_H^{\Delta T}$$

- $g \in \Lambda_H^{\Delta T}$ is defined as $\langle g, \mu \rangle_{\Gamma T} := b_\Gamma(\bar{\mathbf{u}}_h^{\Delta t}, \mu) \quad \forall \mu \in \Lambda_H^{\Delta T}$

Lemma (Operator form)

Equivalent operator form is: find $\lambda_H^{\Delta T} \in \Lambda_H^{\Delta T}$ such that

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Space-time Steklov–Poincaré operator

Lemma (Equivalence)

The MMMFE method is **equivalent to**: find $\lambda_H^{\Delta T} \in \Lambda_H^{\Delta T}$ such that

$$-b_\Gamma^T(\mathbf{u}_h^{\Delta t,*}(\lambda_H^{\Delta T}), \mu) = b_\Gamma^T(\bar{\mathbf{u}}_h^{\Delta t}, \mu) \quad \forall \mu \in \Lambda_H^{\Delta T}.$$

Space-time Steklov–Poincaré operator

- $S : \Lambda_H^{\Delta T} \rightarrow \Lambda_H^{\Delta T}$

$$\langle S\lambda, \mu \rangle_{\Gamma T} := \sum_i \langle S_i \lambda, \mu \rangle_{\Gamma_i^T}, \quad \langle S_i \lambda, \mu \rangle_{\Gamma_i^T} := -\langle \mathbf{u}_h^{\Delta t,*}(\lambda) \cdot \mathbf{n}_i, \mu \rangle_{\Gamma_i^T} \quad \forall \lambda, \mu \in \Lambda_H^{\Delta T}$$

- $g \in \Lambda_H^{\Delta T}$ is defined as $\langle g, \mu \rangle_{\Gamma T} := b_\Gamma^T(\bar{\mathbf{u}}_h^{\Delta t}, \mu) \quad \forall \mu \in \Lambda_H^{\Delta T}$

Lemma (Operator form)

Equivalent operator form is: find $\lambda_H^{\Delta T} \in \Lambda_H^{\Delta T}$ such that

$$S \lambda_H^{\Delta T} = g.$$

Spectral bound, space-time domain decomposition algorithm

Theorem (Spectral bound)

Let the mortar assumptions hold. Then the operator S is **positive definite**. Let moreover $\mathcal{T}_{h,i}$ be **quasi-uniform** and $h \leq Ch_i$ for all i . Then the following **spectral bound** holds:

$$\forall \mu \in \Lambda_H^{\Delta T}, \quad C_0 \|\mu\|_{\Gamma T}^2 \leq \langle S\mu, \mu \rangle_{\Gamma T} \leq C_1 h^{-1} \|\mu\|_{\Gamma T}^2.$$

Comments

- well-posed space-time interface problem
- leads to a **space-time domain decomposition** algorithm
- GMRES can be applied; convergence through the field-of-values estimates:

$$\|\mathbf{r}_k\| \leq \left(\sqrt{1 - (C_0/C_1)^2 h^2} \right)^k \|\mathbf{r}_0\|$$

- on all iterations: problems posed in the **individual space-time subdomains** Ω_i^T and **solved in parallel** (no synchronization on time steps)

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Outline

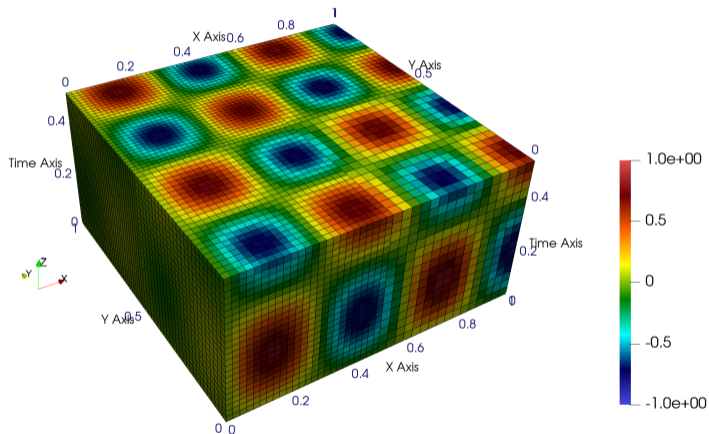
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Numerical experiments

Setting

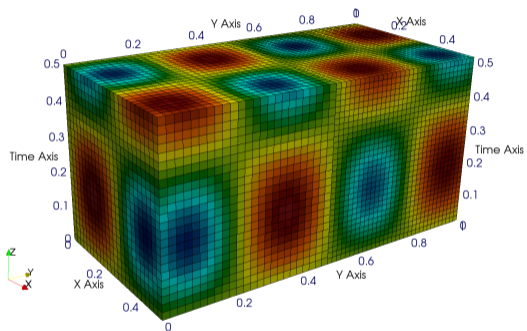
- $d = 2$
- $\mathbf{V}_{h,i} \times W_{h,i}$ on each Ω_i is the lowest-order Raviart–Thomas pair $RT_0 \times DGQ_0$ ($k = l = 0$)
- backward Euler time discretization in each Ω_i^T ($q = 0$)
- mortar finite element space $\Lambda_{H,ij}^{\Delta T}$: discontinuous bilinear ($m = s = 1$, $H = 2h$ and $\Delta T = 2\Delta t$) and discontinuous biquadratic ($m = s = 2$, $H = \sqrt{h}$ and $\Delta T = \sqrt{\Delta t}$) mortars
- GMRES without preconditioner for the space-time interface problem
- deal.II package
- $\Delta t^{-\frac{1}{2}}$ loss in convergence rate in the theoretical bound not observed in the numerical results

Example 1, smooth solution, Ω^T

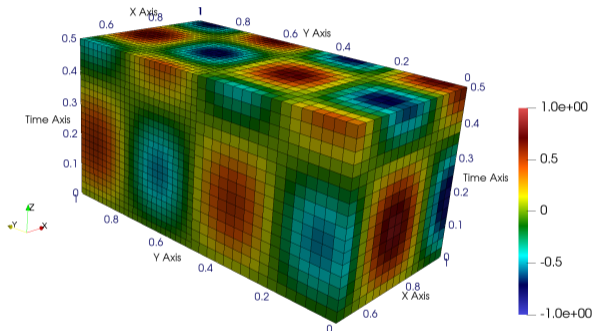


Pressure, bilinear mortars $m = s = 1$, space-time grid at refinement 3, whole Ω^T

Example 1

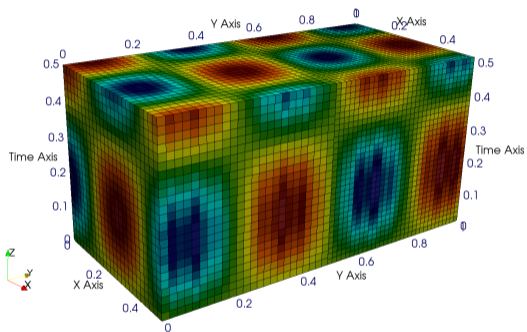


Pressure detail on $\Omega_1^T \cup \Omega_4^T$

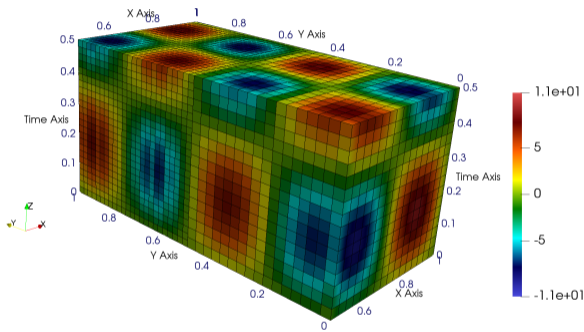


Pressure detail on $\Omega_2^T \cup \Omega_3^T$

Example 1



x-velocity detail on $\Omega_1^T \cup \Omega_4^T$



x-velocity detail on $\Omega_2^T \cup \Omega_3^T$

Example 1

Ref. No.	Ω_1^T			Ω_2^T			Ω_3^T			Ω_4^T			$\Gamma^T(m=1)$			$\Gamma^T(m=2)$		
	n_1	N_1	#DoF	n_2	N_2	#DoF	n_3	N_3	#DoF	n_4	N_4	#DoF	n_Γ	N_Γ	#DoF	n_Γ	N_Γ	#DoF
0	3	3	33	2	2	16	4	4	56	3	3	33	1	1	16	1	1	36
1	6	6	120	4	4	56	8	8	208	6	6	120	2	2	64			
2	12	12	456	8	8	208	16	16	800	12	12	456	4	4	256	2	2	144
3	24	24	1776	16	16	800	32	32	3136	24	24	1776	8	8	1024			
4	48	48	7008	32	32	3136	64	64	12416	48	48	7008	16	16	4096	4	4	576

Meshes, polynomial degrees, and number of degrees of freedom

Example 1: $m = s = 2$ better than $m = s = 1$

Ref.	# GMRES		$\ \mathbf{u} - \mathbf{u}_h^{\Delta t}\ _{L^2(0,T;L^2(\Omega))}$		$\ p - p_h^{\Delta t}\ _{DG}$		$\ p - p_h^{\Delta t}\ _{L^2(0,T;W)}$		$\ \lambda - \lambda_H^{\Delta T}\ _{L^2(0,T;\Lambda_H)}$	
0	11	Rate	6.50e-01	Rate	1.21e+00	Rate	7.91e-01	Rate	7.98e-01	Rate
1	23	-1.06	3.63e-01	0.84	7.21e-01	0.75	4.76e-01	0.73	5.11e-01	0.64
2	39	-0.76	1.74e-01	1.06	3.19e-01	1.18	2.46e-01	0.95	2.34e-01	1.13
3	59	-0.60	8.63e-02	1.02	1.46e-01	1.13	1.25e-01	0.98	1.20e-01	0.96
4	86	-0.54	4.29e-02	1.01	6.93e-02	1.08	6.25e-02	1.00	6.11e-02	0.97

Convergence with bilinear mortars $m = s = 1$

Ref.	# GMRES		$\ \mathbf{u} - \mathbf{u}_h^{\Delta t}\ _{L^2(0,T;L^2(\Omega))}$		$\ p - p_h^{\Delta t}\ _{DG}$		$\ p - p_h^{\Delta t}\ _{L^2(0,T;W)}$		$\ \lambda - \lambda_H^{\Delta T}\ _{L^2(0,T;\Lambda_H)}$	
0	18	Rate	6.81e-01	Rate	1.35e+00	Rate	8.39e-01	Rate	2.13e+00	Rate
2	34	-0.46	1.70e-01	1.00	3.51e-01	0.97	2.51e-01	0.87	2.82e-01	1.46
4	57	-0.37	4.48e-02	0.96	8.59e-02	1.02	6.59e-02	0.96	9.20e-02	0.81

Convergence with biquadratic mortars $m = s = 2$

Example 1: $m = s = 2$ better than $m = s = 1$

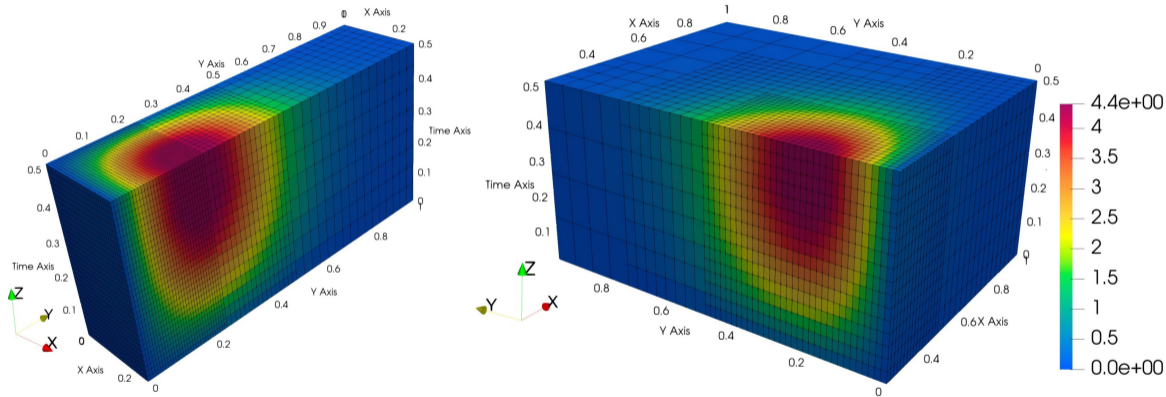
Ref.	# GMRES		$\ \mathbf{u} - \mathbf{u}_h^{\Delta t}\ _{L^2(0,T;L^2(\Omega))}$		$\ p - p_h^{\Delta t}\ _{DG}$		$\ p - p_h^{\Delta t}\ _{L^2(0,T;W)}$		$\ \lambda - \lambda_H^{\Delta T}\ _{L^2(0,T;\Lambda_H)}$	
0	11	Rate	6.50e-01	Rate	1.21e+00	Rate	7.91e-01	Rate	7.98e-01	Rate
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Ref.	# GMRES		$\ \mathbf{u} - \mathbf{u}_h^{\Delta t}\ _{L^2(0,T;L^2(\Omega))}$		$\ p - p_h^{\Delta t}\ _{DG}$		$\ p - p_h^{\Delta t}\ _{L^2(0,T;W)}$		$\ \lambda - \lambda_H^{\Delta T}\ _{L^2(0,T;\Lambda_H)}$	
0	18	Rate	6.81e-01	Rate	1.35e+00	Rate	8.39e-01	Rate	2.13e+00	Rate
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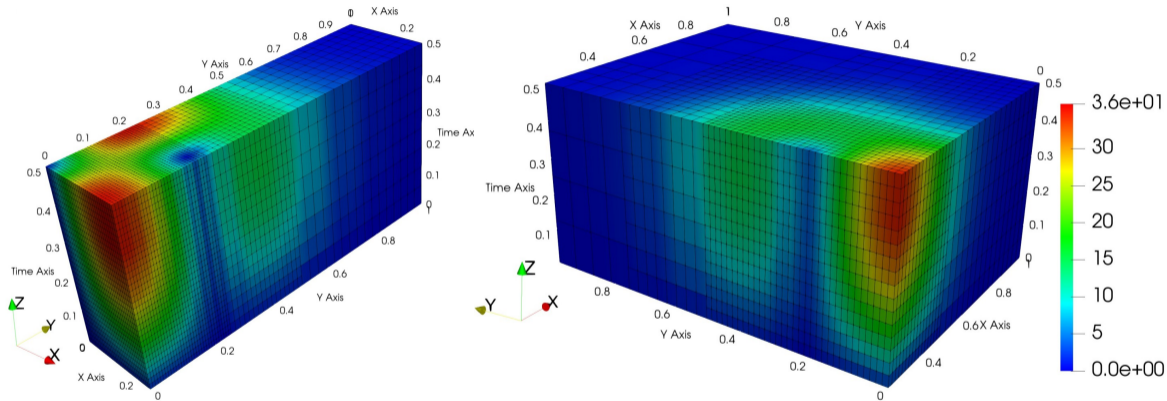
Convergence with biquadratic mortars $m = s = 2$

Example 2, sharp boundary layer, discontinuous bilinear mortars ($m = s = 1$)



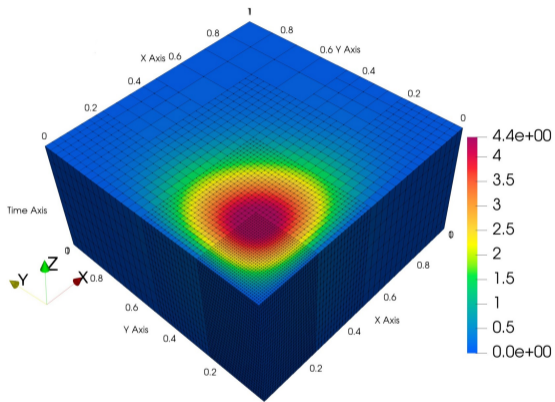
Pressure, cut along the plane $x = 0.25$

Example 2

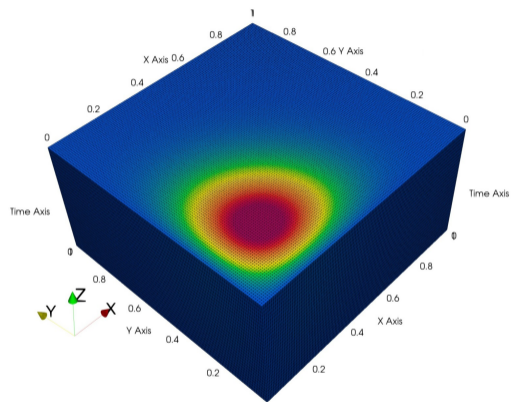


Velocity magnitude, cut along the plane $x = 0.25$

Example 2, Ω^T

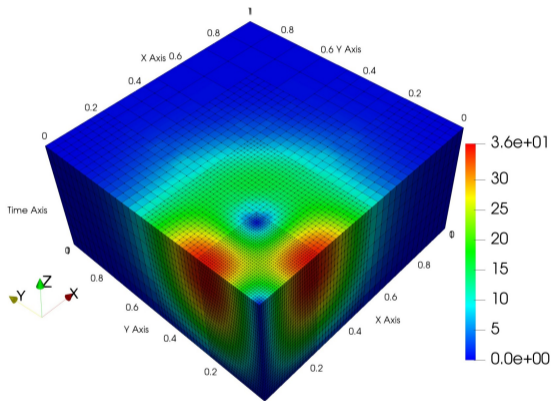


Pressure, mortar multiscale method

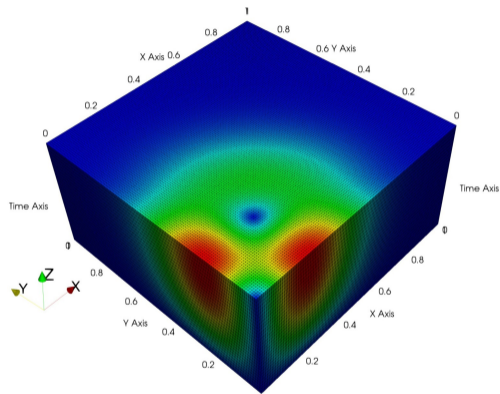


Pressure, fine-scale method

Example 2, Ω^T



Velocity magnitude, mortar multiscale method



Velocity magnitude, fine-scale method

Example 2

Method	# GMRES	$\ \mathbf{u} - \mathbf{u}_h^{\Delta t}\ _{L^2(0,T;L^2(\Omega))}$	$\ p - p_h^{\Delta t}\ _{DG}$	$\ p - p_h^{\Delta t}\ _{L^2(0,T;W)}$	$\ \lambda - \lambda_H^{\Delta T}\ _{L^2(0,T;\Lambda_H)}$
multiscale	102	5.657e-02	8.425e-02	6.319e-02	5.796e-02
fine-scale	140	1.524e-02	2.234e-02	2.154e-02	3.016e-02

Errors and GMRES iterations for the multiscale and fine-scale methods

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- mortar coupling: **space-time interface problem**
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Future directions

- developing a preconditioner for the space-time interface iterative solver
- dedicated a posteriori error analysis

References

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Thank you for your attention!

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