Greedy reconstruction algorithm for the identification of spin distribution

S. Buchwald,¹ G. Ciaramella^{, 2} J. Salomon,³ and D. Sugny^{4,*}

¹Universität Konstanz, Universitätsstraße 10, 78464 Konstanz, Germany

²MOX-Laboratory, Dipartimento di Matematica, Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133 Milano, Italy

³INRIA Paris, ANGE team, Building A, Office 323, 2 rue Simone Iff, 75589 Paris, France

⁴Laboratoire Interdisciplinaire Carnot de Bourgogne (ICB), UMR 6303 CNRS–Université Bourgogne-Franche Comté,

9 Av. A. Savary, B.P. 47 870, F-21078 Dijon Cedex, France



(Received 26 August 2021; accepted 1 December 2021; published 13 December 2021)

We propose a greedy reconstruction algorithm to find the probability distribution of a parameter characterizing an inhomogeneous spin ensemble in nuclear magnetic resonance. The identification is based on the application of a number of constant control processes during a given time for which the final ensemble magnetization vector is measured. From these experimental data, we show that the identifiability of a piecewise constant approximation of the probability distribution is related to the invertibility of a matrix, which depends on the different control protocols applied to the system. The algorithm aims to design specific controls which ensure that this matrix is as far as possible from a singular matrix. Numerical simulations reveal the efficiency of this algorithm in different examples. A systematic comparison with respect to random constant pulses is done.

DOI: 10.1103/PhysRevA.104.063112

I. INTRODUCTION

The identification of parameters that characterize the dynamics of a quantum system is a fundamental prerequisite for controlling its evolution [1-8] and is of practical interest for realizing specific tasks in quantum technologies [9]. This aspect is crucial in open-loop configurations for which the control protocols are designed without any experimental feedback from the system during the control process [1,2,5,10,11]. In the context of quantum systems, the problem of identifying unknown parameters (or functions) has been explored in a large number of studies and for a variety of applications ranging from molecular physics [12-14] and magnetic resonance [15–18] to quantum information science [19–29] and open quantum systems [30-33]. Some mathematical results have also been established in this direction [34-42]. On the basis of different measurement processes and specific control protocols, the goal of these works is generally to estimate the value of one or several parameters of the system Hamiltonian. When controlling an ensemble of identical quantum systems, such a parameter may vary in a given range due to experimental limitations or uncertainties. A key example comes from the spatial inhomogeneities of the external control [2,43-47]. In this case, all the systems are not subjected to exactly the same control. This aspect has to be taken into account in the modeling of the dynamics and in the computation of the control procedure. Robust control protocols against such inaccuracies have been developed recently to solve this experimental issue [2,45,48–55]. However, the variation range of the unknown parameter is not the only crucial quantity; the probability distribution of this parameter (i.e., the number of systems for

each value of the parameter) may also play a major role. It is generally assumed that this probability distribution is flat or has a simple Gaussian or Lorentzian form. In these cases, the probability distribution can be quite easily characterized. However, the problem of identifying probability distributions becomes much more difficult if these have complex structures, with, e.g., several peaks, or if no information is known about them. It is therefore essential to be able to identify with a great precision these unknown probability distributions.

This paper aims at taking a step toward the answer to this open question by developing a numerical algorithm, called a greedy reconstruction algorithm (GRA). By definition, an algorithm is said to be greedy if it takes the best choice available at each iterative step. Greedy algorithms generally find a suboptimal solution, but in a computational time which may be very small compared to that of a global optimization procedure. Such algorithms have recently been applied to the identification of quantum systems [56,57] and we propose to adapt them to the reconstruction of the probability distribution. For the sake of clarity, we focus in this study on a specific example, although our algorithm applies to a large variety of systems. We consider the case of a spin ensemble in nuclear magnetic resonance (NMR) [44,58–61] subjected to an inhomogeneous radio-frequency magnetic field whose range of variation is known, but not its probability distribution. The probability distribution is approximated by a piecewise constant function taking at most K values. The algorithm then designs a series of K controls for the GRA (or less for the optimized version) which are, in the second step, applied to the spin ensemble. The K measured ensemble magnetization vectors at the final time are then used to identify the probability distribution. More precisely, the identification process is related to the invertibility of a matrix, which depends on the different controls. The aim of the algorithm is therefore

^{*}dominique.sugny@u-bourgogne.fr

to design specific control protocols which ensure that this matrix is as far as possible from being singular. The precision of the identification process can be understood from the eigenvalues and eigenvectors of the matrix. In the examples analyzed in this study, constant controls will be sufficient to find the probability distribution with a very good accuracy, but time-dependent controls could also be used. We show that the optimization procedure of the algorithm has a unique solution and a good convexity structure leading to fast convergence. We point out that the controls depend only on the model system and not on the spin distribution or on the available data. The numerical efficiency of the algorithms is shown in different illustrative examples, namely, a double-peak distri-

constant pulses is also done. The paper is organized as follows. Section II describes the model system. Section III is dedicated to the theoretical framework, while the algorithm is presented in Sec. IV. The efficiency of the GRA is numerically demonstrated in Sec. V in two standard examples, a double-peak and a step probability distribution. The conclusion and prospective views are given in Sec. VI. A mathematical description of the method and the development of an optimized greedy algorithm are reported in Appendixes A and B, respectively. Additional numerical results are presented in the Supplemental Material [62].

bution and a step one. A systematic comparison with random

II. THE MODEL SYSTEM

To illustrate our study, we consider a basic control problem in NMR, i.e., a spin ensemble subjected to an inhomogeneous radio-frequency magnetic field [43–46,61]. In a given rotating frame, we assume that all the spins have the same resonance offset ω . Each isochromat is characterized by a Bloch vector, $\mathbf{M}(\alpha) = [M_x, M_y, M_z]^\mathsf{T}$, whose dynamics are governed by the equations of motion

$$\dot{M}_x = -\omega M_y + (1+\alpha)\omega_y M_z,$$

$$\dot{M}_y = \omega M_x - (1+\alpha)\omega_x M_z,$$

$$\dot{M}_z = (1+\alpha)\omega_x M_y - (1+\alpha)\omega_y M_x,$$

where the coordinates of the Bloch vector satisfy $M_x^2 + M_y^2 + M_z^2 = M_0^2$, with M_0 the equilibrium magnetization. ω_x and ω_y are time-dependent controls that correspond to the components of the magnetic field along the x and the y directions. In this study, we assume that these controls are constant in time. We show in Sec. V that this hypothesis is sufficient for the different examples to identify the probability distributions. The parameter α is used to model the control field inhomogeneities, which are of the order of a few percent in standard experiments [60]. The controls $\frac{\omega_x}{2\pi}$ and $\frac{\omega_y}{2\pi}$ are expressed in Hz. We consider a typical field amplitude ω_0 that can be fixed, for instance, at $\omega_0 = 2\pi \times 100$ Hz. We introduce normalized coordinates as follows:

$$u_x = 2\pi \frac{\omega_x}{\omega_0}; \quad u_y = 2\pi \frac{\omega_y}{\omega_0}; \quad t' = \frac{\omega_0}{2\pi}t^2$$

 $\Delta = 2\pi \frac{\omega}{\omega_0}; \quad \mathbf{X} = \frac{\mathbf{M}}{M_0}.$

We omit the prime in the time below to simplify the notation. We deduce that the differential system can be expressed in normalized units as

$$\begin{aligned} \dot{x} &= -\Delta y + (1+\alpha)u_y z, \\ \dot{y} &= \Delta x - (1+\alpha)u_x z, \\ \dot{z} &= (1+\alpha)u_x y - (1+\alpha)u_y x, \end{aligned} \tag{1}$$

with $x^2 + y^2 + z^2 = 1$. The initial state of the dynamics for each spin is the thermal equilibrium point, i.e., the north pole of the Bloch sphere, $\mathbf{X}_0 = (0, 0, 1)^{\mathsf{T}}$. We neglect the relaxation effect and we consider a control time of the order of 100 ms. This corresponds to a normalized time t_f of the order of 10. In the numerical simulations, we add the constraints $|u_x| \leq u_m$ and $|u_y| \leq u_m$, where u_m is the maximum amplitude of each component. In NMR, only the first two coordinates of the magnetization vector can be directly measured. We do not have access to the *z* component due to the strong magnetic field applied along this direction. We denote by $\mathbf{Y}(t) = (x(t), y(t))^{\mathsf{T}}$ the projection of the Bloch vector onto the first two coordinates. We point out that this aspect is not a limiting point for the application of the identification process.

III. IDENTIFICATION OF THE SPIN DISTRIBUTION

We consider an ensemble of *N* spins whose dynamics are governed by Eq. (1). We assume that the control amplitudes (u_x, u_y) belong to the admissible set $\mathcal{U} = \{(u_x, u_y) \in \mathbb{R}^2 \mid |u_x| \leq u_m, |u_y| \leq u_m\}$. The objective of the control procedure is to identify the probability distribution of the parameter α . To simplify the recognition process, we assume that the ensemble of spins can be decomposed into a set of *K* subgroups with the same value of the parameter α_ℓ , $1 \leq \ell \leq K$. However, the discrete probability distribution P_{\star} for α , namely, the probability of each possible outcome α_ℓ or, in other words, the number of elements N_ℓ in each subgroup, is unknown. By

definition, we have
$$P_{\star}(\ell) = \frac{N_{\ell}}{N}$$
 and $\sum_{\ell=1}^{K} P_{\star}(\ell) = 1$.

The projected solution onto the first two coordinates of Eq. (1) at time t_f is denoted $\mathbf{Y}_{\mathbf{u},\alpha}(t_f)$, where the dependence on \mathbf{u} and α has been explicitly mentioned. The corresponding experimental realization of this controlled dynamic leads to $\mathbf{Y}_{\mathbf{u}}^{\exp}(t_f) = (x_{\mathbf{u}}^{\exp}(t_f), y_{\mathbf{u}}^{\exp}(t_f))^{\mathsf{T}}$, where $\mathbf{Y}_{\mathbf{u}}^{\exp}(t_f)$ can be viewed as the average at time t_f of the experimental measures of all the spins in the set subjected to the control \mathbf{u} . The coordinates $x_{\mathbf{u}}^{\exp}$ and $y_{\mathbf{u}}^{\exp}$ are those of this measured magnetization vector.

The relation between the theoretical description of the dynamical system and the experimental outcome can be expressed as

$$\mathbf{Y}_{\mathbf{u}}^{\exp}(t_f) = \sum_{\ell=1}^{K} P_{\star}(\ell) \mathbf{Y}_{\mathbf{u},\alpha_{\ell}}(t_f), \qquad (2)$$

in which the two sides of the equation crucially depend on the control **u**. A specific control protocol is not sufficient to identify the probability distribution P_{\star} , which generally requires the implementation of *K* control processes with *K* different controls denoted \mathbf{u}_k , k = 1, ..., K. Note that in the optimized version of the GRA presented in Appendix B, the number of controls can be different from *K*. On the basis of the experimental outputs, a straightforward way to determine P_* is to solve the minimization problem

$$\min_{P \in \mathbb{P}} \sum_{k=1}^{K} \left\| \mathbf{Y}_{\mathbf{u}_{k}}^{\exp}(t_{f}) - \sum_{\ell=1}^{K} P(\ell) \mathbf{Y}_{\mathbf{u}_{k},\alpha_{\ell}}(t_{f}) \right\|^{2}, \quad (3)$$

where \mathbb{P} is the set of all the possible probability distribu-

tions *P* that satisfy
$$P(\ell) \ge 0$$
 for $1 \le \ell \le K$ and $\sum_{\ell=1}^{K} P(\ell) = 1$
Mathematically, we point out that \mathbb{P} is a convex and closed

Mathematically, we point out that \mathbb{P} is a convex and closed set. $\|\cdot\|$ denotes the standard Euclidean vector norm. Note that Eq. (3) can be rewritten as

$$\min_{P \in \mathbb{P}} \sum_{k=1}^{K} \left\| \sum_{\ell=1}^{K} (P_{\star}(\ell) - P(\ell)) \mathbf{Y}_{\mathbf{u}_{k},\alpha_{\ell}}(t_{f}) \right\|^{2}.$$
 (4)

At this point, it is clear that a key ingredient of the accuracy of the identification process relies on the choice of the controls \mathbf{u}_k .

To clarify this problem, we introduce a set $\{\varphi_j\}_{j=1}^K$ of linearly independent functions $\varphi_j : \{1, \ldots, K\} \to \mathbb{R}$ such that $\mathbb{P} \subset \text{span}(\{\varphi_j\}_{j=1}^K)$, where span denotes the vector space generated by the functions. Expressing P_{\star} and P as, respectively,

$$P_{\star}(\ell) = \sum_{j=1}^{n} \beta_{\star,j} \varphi_j(\ell)$$
 and $P(\ell) = \sum_{j=1}^{n} \beta_j \varphi_j(\ell)$, the minimiza-

tion problem, (4), becomes

$$\min_{\boldsymbol{\beta}\in\widehat{\mathbb{R}}^{K}}\sum_{k=1}^{K}\left\|\sum_{\ell,j=1}^{K}(\boldsymbol{\beta}_{\star,j}-\boldsymbol{\beta}_{j})\varphi_{j}(\ell)\mathbf{Y}_{\mathbf{u}_{k},\boldsymbol{\alpha}_{\ell}}(t_{f})\right\|^{2},\qquad(5)$$

where the vector $\beta = (\beta_j)_{j=1}^K$ is taken in $\widehat{\mathbb{R}}^K$, a subset of \mathbb{R}^K , so that $P = \sum_j \beta_j \varphi_j$ is a probability distribution. Note that

other basis functions can be used, such as, e.g., a basis of polynomial functions, which may be more suited to specific problems. Equation (5) can be rewritten in a compact form as

$$\min_{\beta \in \widehat{\mathbb{R}}^{K}} \langle \beta_{\star} - \beta | W | \beta_{\star} - \beta \rangle, \tag{6}$$

where *W* is a symmetric and positive semidefinite $K \times K$ matrix whose elements are defined as

$$W_{\ell,j} = \sum_{k} \langle \gamma_{\ell}(\mathbf{u}_{k}) | \gamma_{j}(\mathbf{u}_{k}) \rangle \tag{7}$$

with

$$\gamma_j(\mathbf{u}_k) = \sum_{\ell} \varphi_j(\ell) \mathbf{Y}_{\mathbf{u}_k, \alpha_\ell}(t_f)$$

Since the set of vectors β is a convex subset of \mathbb{R}^{K} , we deduce that the problem is uniquely solvable if the matrix W is positive definite, i.e., if W has a nonzero determinant. In the case where W has a nontrivial kernel, infinitely many solutions may exist which lead to incorrect probability distributions different from the experimental one P_{\star} . We stress that the nontriviality of the kernel depends completely on the choice of the controls \mathbf{u}_{k} .

We show in this study that the GRA allows us to design a set of controls \mathbf{u}_k so that the matrix W is positive definite with

a trivial kernel. Note that larger times and/or larger bounds for the control pulses generally ensure a better conditioning of matrix W and, hence, a better identifiability of the distribution.

The algorithm is composed of two steps, namely, an offline and an online phase. In the first phase (off-line), the GRA computes the controls \mathbf{u}_k . In this phase, only the theoretical model is needed, without any experimental input. The derived controls are used in the second phase (online), in which the different magnetization vectors are measured and the minimization problem, (3), is solved. We thus stress that the experimental data are used only in the online phase of the algorithm. Note that the controls are the same for any probability distribution to identify and depend only on the model system under study. Finally, we point out that, while in the first algorithm we consider that all control pulses have the same duration t_f , in the second version, described in Sec. IV B, the duration of each pulse is considered as a variable to be optimized together with its amplitude. The generality of the GRA allows one to tackle this situation in a straightforward manner.

IV. THE GREEDY RECONSTRUCTION ALGORITHM

We present in this section the GRA in its classical form; an optimized extension called the optimized GRA (OGRA) is described in Appendix B. For pedagogical purposes, we have limited the mathematical derivation of the algorithm to its strict minimum. The interested reader can find mathematical details about the algorithms in [56] and [57] for the standard and optimized GRA, respectively.

A. Optimizing the control amplitudes for a fixed control time

The GRA computes the controls \mathbf{u}_k by solving a sequence of fitting-step and discriminatory-step problems, in which the goal of the first step is to identify a nontrivial kernel of a submatrix of W, while the second phase designs a new control which is aimed at correcting this discrepancy and eliminating the identified nontrivial kernel. The explicit formulation of the algorithm is given in terms of the function $\mathbf{h}^{(k)}$, defined by

$$\mathbf{h}^{(k)}(\boldsymbol{\beta}, \mathbf{u}) = \sum_{\ell=1}^{K} \sum_{j=1}^{k} \beta_{j} \varphi_{j}(\ell) \mathbf{Y}_{\mathbf{u}, \alpha_{\ell}}(t_{f})$$
(8)

for any β in \mathbb{R}^k . The GRA is described below. Some mathematical statements of the different steps of the algorithm are described in Appendix A. Its numerical implementation is presented and discussed in Sec. V.

Greedy reconstruction algorithm

Given a set of *K* linearly independent functions $(\varphi_1, \ldots, \varphi_K)$.

Solve the initialization problem

$$\max_{\mathbf{u}\in\mathcal{U}}\|\mathbf{h}^{(1)}(1,\mathbf{u})\|^2,\tag{9}$$

which gives the control \mathbf{u}_1 , and set k = 1. While $k \leq K - 1$ (1) <u>Fitting step</u>: Find $(\beta_i^k)_{j=1,\dots,k}$ that solves the problem

$$\min_{\boldsymbol{\beta}\in\mathbb{R}^k}\sum_{m=1}^k \|\mathbf{h}^{(K)}(\mathbf{e}_{k+1},\mathbf{u}_m)-\mathbf{h}^{(k)}(\boldsymbol{\beta},\mathbf{u}_m)\|^2, \qquad (10)$$

where \mathbf{e}_{k+1} is the (k+1)th canonical vector in \mathbb{R}^{K} .

(2) <u>Discriminatory step</u>: Find \mathbf{u}_{k+1} that solves the problem

$$\max_{\mathbf{u}\in\mathcal{U}}\|\mathbf{h}^{(K)}(\mathbf{e}_{k+1},\mathbf{u})-\mathbf{h}^{(k)}(\beta_k,\mathbf{u})\|^2.$$
(11)

(3) Update $k + 1 \rightarrow k$.

End while

The basic principles of the GRA can be detailed by its first two iterations for K = 2. Using

$$\mathbf{h}^{(1)}(1,\mathbf{u}) = \sum_{\ell=1}^{K} \varphi_1(\ell) \mathbf{Y}_{\mathbf{u},\alpha_\ell}(t_f)$$

the initialization problem can be expressed as

$$\max_{\mathbf{u}\in\mathcal{U}}\|W_{11}(\mathbf{u})\|^2.$$

We deduce that the goal of this step is to maximize the modulus of this W-matrix element, so as to be as far as possible from a 0 of W. We then consider the first step of the algorithm with k = 1. For clarity we omit below the dependence of W on **u**. By definition, we have

$$W_{11} = \left\| \sum_{\ell=1}^{K} \varphi_1(\ell) \mathbf{Y}_{\mathbf{u},\alpha_{\ell}}(t_f) \right\|^2,$$

$$W_{22} = \left\| \sum_{\ell=1}^{K} \varphi_2(\ell) \mathbf{Y}_{\mathbf{u},\alpha_{\ell}}(t_f) \right\|^2,$$

$$W_{12} = W_{21} = \sum_{\ell,\ell'=1}^{K} \varphi_1(\ell) \varphi_2(\ell') \langle \mathbf{Y}_{\mathbf{u},\alpha_{\ell}}(t_f) | \mathbf{Y}_{\mathbf{u},\alpha_{\ell'}}(t_f) \rangle,$$

and we deduce that the quantity to minimize in the fitting step can be written as

$$\|\mathbf{h}^{(2)}(e_2,\mathbf{u}_1) - \mathbf{h}^{(1)}(\beta,\mathbf{u}_1)\|^2 = W_{11}\beta^2 - 2W_{12}\beta + W_{22},$$

where β is here a real number. The minimum is reached for $\beta_1 = W_{11}^{-1}W_{12}$, where *W* is computed for the control **u**₁. This value can be associated with a vector $(\beta_1, -1)^{\mathsf{T}}$ of the kernel of the following 2 × 2 submatrix of *W*:

$$\begin{pmatrix} W_{11} & W_{12} \\ W_{12} & W_{22} \end{pmatrix}.$$
 (12)

The fitting step of the GRA can thus be interpreted as a systematic way to find a basis of the kernel of larger and larger submatrices of W. Setting β to β_1 , the discriminatory step consists in adding a new control \mathbf{u}_2 to correct this singularity, i.e., in selecting this control such that the corresponding quantity is as far as possible from a 0. Mathematically, it can be shown that this procedure always has a solution and that the new matrix $W^{(2)}$ has a nontrivial 2×2 submatrix, (12) (see Appendix A for details).





FIG. 1. Plot of the true distribution (red crosses, which essentially coincide with the back circles) and approximated solutions with minimum error as a function of α , computed by the different optimization algorithms for the identification problem using control sets with a fixed control time. In particular, controls generated by the OGRA (green lines), GRA (filled black circles), and RCC (blue squares) are plotted. Dimensionless units are used.

B. Optimizing the amplitude controls and time horizon

Thus far, we have considered a fixed control time t_f . However, it is also possible to consider controls with different control times, up to a fixed boundary $t_f^{\max} > 0$. In this case, we also maximize with respect to time, meaning that the initialization and discriminatory step problems at iteration kwould change to

$$\max_{\substack{\mathbf{u}\in\mathcal{U},\\t_f\in[0,t_f^{\max}]}} \|\mathbf{h}^{(1)}(1,\mathbf{u};t_f)\|^2$$
(13)

and

10000

$$\max_{\substack{\mathbf{u}\in\mathcal{U},\\t_f\in[0,f^{\max}]}} \|\mathbf{h}^{(K)}(e^{k+1},\mathbf{u};t_f) - \mathbf{h}^{(k)}(\beta^k,\mathbf{u};t_f)\|^2, \qquad (14)$$

respectively. In Eqs. (13) and (14), the function $\mathbf{h}^{(k)}$ is still defined as in (8), only with the control time as an additional variable. Similarly, one can adapt the corresponding problems in the OGRA, the optimized version described in Appendix B. We denote the resulting two algorithms GRAt and OGRAt.

V. NUMERICAL RESULTS

A. The case of a double-peak distribution

As the first illustrative example, we investigate here the identification of a symmetric double-peak probability distribution, displayed in Fig. 1. Similar results have been achieved for other smooth distributions with one or several peaks. Numerical details are described in the Supplemental Material [62].

In the numerical simulations, we consider a control time $t_f = t_f^{\text{max}} = 16$. The amplitude u_m is equal to 10. The normalized offset resonance is set for all spins to $\frac{\pi}{10}$, i.e., to 30 Hz. We also assume that $\alpha \in [-0.2, 0.2]$ and K = 30. The K discrete

TABLE I. Minimum relative norm error for different control sets.

| | | Control set | | | | | | |
|------------|--------|-------------|--------|--------|--------|--------|--|--|
| | GRA | GRAt | OGRA | OGRAt | RCC | RCCt | | |
| Min. error | 0.0045 | 0.0098 | 0.0005 | 0.0009 | 0.4685 | 0.0841 | | |

values α_{ℓ} of α are regularly spaced in the interval of variation of α , i.e., $\alpha_{\ell} = -0.2 + 0.4 \frac{\ell-1}{K-1}$. Since the control protocols are constant in time, Eq. (1) is solved numerically by directly evaluating the exponential matrix corresponding to the exact solution. All optimization problems are solved by a BFGS descent-direction method. We also mention that the exact number of uncoupled spins in the ensemble is not relevant for all the computations, since we are only interested in their probability distribution. However, we use a total number of 10^5 spins in the numerical simulations.

For the GRA and GRAt, we consider a random and orthonormal basis, $\{\varphi_k\}_{k=1}^{30}$. Note that any basis of this space can be used in the respective algorithms. For the OGRA and OGRAt, we extend the basis from the GRA by 30 randomly chosen probability distributions $\{\varphi_k\}_{k=31}^{60}$. The tolerance used in the OGRA and OGRAt (see Appendix B) is set to be tol = 10^{-14} . The controls generated by the algorithms and corresponding to the numerical results discussed below are described in the Supplemental Material [62]. To test whether it is even necessary to run the algorithm or whether the same results could be achieved with other control protocols, we also consider two sets of 30 random and constant controls. In the first and second cases, we use completely random constant values in the set \mathcal{U} , with a control time t_f or with different and random control times in the interval $[0, t_f]$, respectively. We denote the two sets of controls RCC and RCCt.

The robustness of the different control functions is evaluated by considering a 30-dimensional hypercube centered in the global minimum P_{\star} of our identification problem, with a radius of $100 \|P_{\star}\|$, and we repeat the minimization process for 100 initialization vectors randomly chosen in this hypercube. We then compute the minimum norm difference $\frac{\|P_{\star}-P_{f}\|}{\|P_{\star}\|}$ over



FIG. 2. Same as Fig. 1, but for variable control times.



FIG. 3. Same as Fig. 1, but for a step distribution.

all optimization runs, where P_f denotes the solution given by the optimization algorithm. We obtain the results reported in Table I.

As can be seen in Table I, the errors of the OGRA and OGRAt are 10 times smaller than those of the GRA and GRAt, which themselves are 10 and 100 times smaller, respectively, than the errors of both sets of random controls. Similar results have been achieved for other smooth distributions, which show the efficiency of the two proposed algorithms. Figures 1 and 2 display the true distribution and the minimal solution for all control sets for fixed and variable control times, respectively.

We observe that the solutions computed with controls generated by any algorithm match the true distribution. On the other hand, the RCC completely fails, showing a third peak in the middle, while the RCCt can identify at least the two peaks of the distribution.

B. The case of a step distribution

As the second illustrative example, we consider a noncontinuous step distribution, displayed in Fig. 3, in which only spins with a positive parameter α can be observed in the sample. We repeat the numerical simulations in Sec. V A and we obtain the results reported in Table II.

As can be seen in Table II, the difference in the magnitude of errors is similar to that for the double-peak distribution. These results are displayed in Figs. 3 and 4.

We observe that the OGRA and OGRAt are still able to identify the true distribution, while the GRA and GRAt already show small discrepancies. The RCC completely fails again, but also the RCCt shows major visible differences in the upper part of the step distribution. Arguments based on

TABLE II. Minimum relative norm error for different control sets.

| | Control set | | | | | | |
|------------|-------------|--------|--------|--------|--------|--------|--|
| | GRA | GRAt | OGRA | OGRAt | RCC | RCCt | |
| Min. error | 0.0295 | 0.0181 | 0.0018 | 0.0021 | 0.4204 | 0.1943 | |



FIG. 4. Same as Fig. 3, but for variable control times.

the properties of matrix W are given in Sec. V C to explain these numerical observations.

C. Eigenvalues and eigenvectors of W

We explain qualitatively in this section the numerical results observed in Secs. V A and V B through the properties of matrix W, i.e., its eigenvalues and eigenvectors. We present the spectra of matrix W for different control sets in Figs. 5 and 6. A very large difference is observed in the eigenvalues associated with the optimized controls versus the random ones. Note that this observation is the same if the control time is also optimized. This difference is quantitatively measured by the condition number of W, i.e., the ratio between the largest and the smallest eigenvalues, which is listed in Table III. As could be expected, these results show that, using random controls, matrix W can be close to being singular. In the example in Fig. 6, while most of the eigenvalues are larger than 1, four of them are smaller than 10^{-10} . Hence, matrix W has a very bad condition number. We stress the very good result achieved



FIG. 5. Spectrum of matrix W defined in (7) for controls generated by the OGRA (dash-dotted green lines), GRA (black circles), and RCC (blue squares). Dimensionless units are used.



FIG. 6. Same as Fig. 5, but for variable control times.

by the OGRA, for which all the eigenvalues have almost the same value. This analysis may also explain the difference between a smooth and a noncontinuous probability distribution. As a matter of fact, numerical results reveal that random controls have more difficulty identifying a nonsmooth probability distribution as illustrated in Sec. V B. This aspect can be understood from the behavior of the eigenvectors. Indeed, we observe numerically that the modes with a large number of oscillations correspond to the smallest eigenvalues. Such modes have to be used to reconstruct probability distributions with rapid and abrupt variations. For random controls, these eigenvectors lead to large errors and to incorrect probability distributions.

VI. CONCLUSIONS

We have introduced in this work a greedy reconstruction algorithm with an application to spin dynamics. The algorithm provides a systematic way to identify the probability distribution of a parameter of a Hamiltonian system varying in a given range. The efficiency of the identification process has been illustrated in the case of a spin ensemble subjected to an inhomogeneous radio-frequency magnetic field. After the description of some mathematical properties of the algorithm, numerical simulations have revealed the efficiency of the GRA and its quite large basin of convergence. We have shown that the GRA is able to identify nontrivial probability distributions with several peaks or with a step variation. An optimized version of this algorithm can be derived to further improve the identification process. We have limited the study to constant controls, but similar results can be achieved with time-dependent pulses. A quantitative comparison with random constant controls has highlighted the nontrivial recognition process realized by the algorithms. The numerical observations can also be partly explained by the computation of the eigenvalues and eigenvectors of matrix W.

This analysis paves the way for further investigations in magnetic resonance. An interesting direction is the study of the sensitivity of the algorithm to experimental imperfections or to the presence of noise. It could also be used to identify the probability distribution of other parameters, such as 4.9×10^{3}

Cond(W)

RCCt 1.42×10^{17}

| | TABLE III. Condition nu | imber of matrix W | for different control s | sets. | |
|-------------|-------------------------|-------------------|-------------------------|-------|--|
| Control set | | | | | |
| GRA | GRAt | OGRA | OGRAt | RCC | |

19.55

the resonance offset. These greedy algorithms could also be transferred to other domains such as quantum optics and atomic and molecular physics. Finally, we hope that our method will be used in relevant experimental applications in magnetic resonance in the near-future. Inhomogeneous contributions may be due, e.g., to specific shapes of the coil used to generate the radio-frequency magnetic field, which leads to spatial variations of the field and thus to a nontrivial probability distribution of the system parameters [63–65].

 6.9×10^{7}

ACKNOWLEDGMENTS

The work of the first author was supported by the DFG via the collaborative research center SFB1432, Project-ID 425217212. This research was partially supported by the ANR project "QUACO," ANR-17-CE40-0007-01. This project received funding from the European Union Horizon 2020 research and innovation program under Marie-Sklodowska-Curie Grant No. 765267 (QUSCO).

APPENDIX A: MATHEMATICAL DESCRIPTION OF THE GRA

We give some mathematical details of the GRA in this Appendix. Straightforward computations show that the different steps of the GRA can be expressed in matrix form as follows:

1. The initialization problem, (9), is equivalent to

$$\max_{\mathbf{u} \in \mathcal{U}} [W(\mathbf{u})]_{1,1}.$$

2. The fitting-step problem, (10), is equivalent to

$$\min_{\beta \in \mathbb{R}^k} \langle \beta | W_{[1:k,1:k]}^k | \beta \rangle - 2 \langle W_{[1:k,k+1]}^k | \beta \rangle,$$

where $W^{k} = \sum_{m=1}^{k} W(u_{m})$. $W^{k}_{[1:k,1:k]}$ and $W^{k}_{[1:k,k+1]}$ denote, re-

spectively, the $k \times k$ upper-left block of W^k and a column vector containing the first k components of the (k + 1)th column of W^k .

3. The discriminatory-step problem, (11), is equivalent to

$$\max_{u\in\mathcal{U}}\langle \mathbf{v}|[W(\mathbf{u})]_{[1:k+1,1:k+1]}|\mathbf{v}\rangle,$$

where $\mathbf{v} = (\beta_k^{\mathsf{T}}, -1)^{\mathsf{T}}$.

The different iterations of the GRA can then be described as follows. At iteration k, we assume that the submatrix $W_{[1:k,1:k]}^k$ is positive definite, but $W_{[1:k+1,1:k+1]}^k$ can have a nontrivial kernel. The idea is first to identify the kernel of $W_{[1:k+1,1:k+1]}^k$ by solving (10) and then to compute a new control \mathbf{u}_{k+1} such that the new updated matrix $W^{k+1} = W^k + W^k$ $W(\mathbf{u}_{k+1})$ has a positive definite upper-left block, $W_{[1:k+1,1:k+1]}^{k+1}$. The convergence of the algorithm follows from this iterative process.

The following two technical lemmas describe the optimizations used in the two steps of the algorithm. In particular, Lemma A shows that the fitting step identifies the kernel of the matrix $W_{[1:k+1,1:k+1]}^k$.

16.6178

 4.3×10^{9}

Lemma 1. Assume that $W_{[1:k,1:k]}^k$ is positive definite and $W_{[1:k+1,1:k+1]}^k$ has a nontrivial kernel. Then the vector $\mathbf{v} =$ $(\beta_k^{\mathsf{T}}, -1)^{\mathsf{T}}$, where β_k is the solution to (10), is in the kernel of $W_{[1:k+1,1:k+1]}^k$.

The second lemma is the basis of the discriminatory-step algorithm and shows that this step corrects the rank deficiency of $W_{[1:k+1,1:k+1]}^k$.

Lemma 2. Let $W_{[1:k,1:k]}^k$ be a positive definite matrix and β_k a solution of the fitting-step problem, (10). Any solution \mathbf{u}_{k+1} of (11) satisfies

$$\langle \mathbf{v} | W_{[1:k+1,1:k+1]} | \mathbf{v} \rangle > 0$$

for k = 0, 1, ..., K - 1, where $\mathbf{v} = (\beta_k^{\mathsf{T}}, -1)^{\mathsf{T}}$.

The mathematical proofs of these results and a detailed numerical analysis of the GRA are beyond the scope of this work and will be presented elsewhere.

APPENDIX B: THE OPTIMIZED GREEDY ALGORITHM

We discuss in this Appendix the optimized version of the GRA. It can be shown numerically that the behavior and the efficiency of the GRA are strongly affected by the choice of the elements φ_k and their ordering. The GRA is essentially a sweep over the set $(\varphi_k)_{k=1}^K$. However, an incorrect choice of the elements φ_k and their ordering can lead to the stagnation of the algorithm and to the computation of many useless control functions. Note that the stagnation of the algorithm can be measured in terms of rank corrections, i.e., whether for consecutive iterations the rank of W does not increase. These reasons are at the origin of an optimized algorithm [57]. The OGRA takes as input a set Φ , possibly larger than $(\varphi_k)_{k=1}^{K}$ with linearly dependent elements, and returns as output not only a set of \widetilde{K} control functions, but also a set of linearly independent functions $(\widehat{\varphi}_k)_{k=1}^{\widehat{K}}$. The integers \widetilde{K} and \widehat{K} are not necessarily equal and may be smaller than K (in contrast to the GRA). The extension of the OGR method in [57] to the distribution reconstruction problem is detailed below, where we use the map $\mathbf{h}_{\mathcal{S}}$, defined as

$$\mathbf{h}_{\mathcal{S}}(\boldsymbol{\beta}, \mathbf{u}) = \sum_{\ell=1}^{K} \sum_{j=1}^{\operatorname{card}[\mathcal{S}]} \boldsymbol{\beta}_{j} \varphi_{j}(\ell) \mathbf{Y}(\mathbf{u}, \alpha_{\ell}),$$

where $S = (\varphi_1, \ldots, \varphi_k)$. Note that for the fitting-step problem, we do not have any constraint for the choice of coefficients β . This is due to the fact that, within the algorithm, we are not trying to reconstruct a distribution but to make the respective submatrix positive definite.

Optimized greedy reconstruction algorithm

Given a set of $K_+ \ge K$ linearly independent functions $(\varphi_1, \ldots, \varphi_{K_+})$ and a tolerance tol > 0.

Solve the initialization problem

$$\max_{n \in \{1, \dots, K_+\}} \max_{\mathbf{u} \in \mathcal{U}} \left\| \mathbf{h}_{\varphi_n}^{(1)}(1, \mathbf{u}) \right\|^2, \tag{B1}$$

which gives the control \mathbf{u}_1 and the control ℓ_1 . Set k = 1 and $S = \{\varphi_{\ell_1}\}, \widetilde{K} = K_+$, and update $\Phi = \Phi \setminus \{\varphi_{\ell_1}\}$. The algorithm is stopped if $\|\mathbf{h}_{S}^{(1)}(1, \mathbf{u})\|^2 < \text{tol.}$

While $k \leq \tilde{K} - 1$ do

(1) Remove elements from Φ that are linearly dependent on the ones in S. Shift the indices of the remaining elements in Φ . Update card $[\Phi] \rightarrow \widetilde{K}$.

(2) For $\ell = 1, \ldots, \tilde{K}$ do

- [1] C. Brif, R. Chakrabarti, and R. Rabitz, New J. Phys. 12, 075008 (2010).
- [2] S. J. Glaser, U. Boscain, T. Calarco, C. P. Koch, W. Köckenberger, R. Kosloff, I. Kuprov, B. Luy, S. Schirmer, T. Schulte-Herbrüggen, D. Sugny, and F. K. Wilhelm, Eur. Phys. J. D 69, 279 (2015).
- [3] C. P. Koch, M. Lemeshko, and D. Sugny, Rev. Mod. Phys. 91, 035005 (2019).
- [4] C. P. Koch, J. Phys.: Condens. Matter 28, 213001 (2016).
- [5] A. Borzì, G. Ciaramella, and M. Sprengel, Formulation and Numerical Solution of Quantum Control Problems (SIAM, Philadelphia, PA, 2017).
- [6] Y. Maday, J. Salomon, and G. Turinici, Numer. Math. 103, 323 (2006).
- [7] M. Sprengel, G. Ciaramella, and A. Borzì, J. Dyn. Control Syst. 24, 657 (2018).
- [8] U. Boscain, M. Sigalotti, and D. Sugny, PRX Quantum 2, 030203 (2021)
- [9] A. Acín, I. Bloch, H. Buhrman, T. Calarco, C. Eichler, J. Eisert, D. Esteve, N. Gisin, S. J. Glaser, F. Jelezko, S. Kuhr, M. Lewenstein, M. F. Riedel, P. O. Schmidt, R. Thew, A. Wallraff, I. Walmsley, and F. K. Wilhelm, New J. Phys. 20, 080201 (2018).
- [10] D. D'Alessandro, Introduction to Quantum Control and Dynamics (Chapman & Hall/CRC, Boca Raton, FL, 2007).
- [11] A. E. Bryson and Y.-C. Ho, *Applied Optimal Control: Optimization, Estimation, and Control* (Hemisphere, Washington, DC, 1975).
- [12] J. M. Geremia and H. Rabitz, J. Chem. Phys. 118, 5369 (2003).
- [13] A. Feng, H. Rabitz, G. Turinici, and C. Le Bris, J. Phys. Chem. A 110, 7755 (2006).
- [14] J. M. Geremia and H. Rabitz, Phys. Rev. Lett. 89, 263902 (2002).
- [15] D. Ma, V. Gulani, and N. Seiberlich, Nature (London) 495, 187 (2013).
- [16] Q. Ansel, M. Tesch, S. J. Glaser, and D. Sugny, Phys. Rev. A 96, 053419 (2017).
- [17] E. Y. Pierre, D. Ma, Y. Chen, C. Badve, and M. A. Griswold, Magn. Reson. Med. 75, 2481 (2016).
- [18] G. Ciaramella, A. Borzì, G. Dirr, and D. Wachsmuth, SIAM J. Sci. Comput. 37, A319 (2015).

<u>Fitting step</u>: Find $(\beta_i^{\ell})_{j=1,\dots,k}$ that solve the problem

$$\min_{\beta \in \mathbb{R}^k} \sum_{m=1}^k \|\mathbf{h}_{\varphi_\ell}(1, \mathbf{u}_m) - \mathbf{h}_{\mathcal{S}}(\beta, \mathbf{u}_m)\|^2, \quad (B2)$$

end for

(3) <u>Discriminatory step</u>: Find \mathbf{u}_{k+1} and ℓ_{k+1} that solve the problem

$$\max_{\ell \in \{1,\dots,K_+\}} \max_{\mathbf{u} \in \mathcal{U}} \|\mathbf{h}_{\varphi_{\ell}}(1,\mathbf{u}) - \mathbf{h}_{\mathcal{S}}(\beta^{\ell},\mathbf{u})\|^2.$$
(B3)

If $\|\mathbf{h}_{\varphi_{\ell}}(1, \mathbf{u}^{k+1}) - \mathbf{h}_{\mathcal{S}}(\beta^{\ell_{k+1}}, \mathbf{u}_{k+1})\|^2 < \text{tol}$, then stop and return \mathcal{S} and the computed $(\mathbf{u})_{m=1}^k$.

(4) Orthogonalize the function $\varphi_{\ell_{k+1}}$ with respect to S and update $S \cup \{\varphi_{\ell_{k+1}}\} \to S$, $\Phi \setminus \{\varphi_{\ell_{k+1}}\} \to \Phi$, and $k + 1 \to k$. **End while**

[19] S. G. Schirmer, A. Kolli, and D. K. L. Oi, Phys. Rev. A 69, 050306(R) (2004).

- [20] S. G. Schirmer and D. K. L. Oi, Phys. Rev. A 80, 022333 (2009).
- [21] S. G. Schirmer and F. C. Langbein, Phys. Rev. A **91**, 022125 (2015).
- [22] H. Yuan and C. F. Fung, npj Quantum Info. 3, 14 (2017).
- [23] J. Liu and H. Yuan, Phys. Rev. A 96, 012117 (2017).
- [24] N. Wittler, F. Roy, K. Pack, M. Werninghaus, A. S. Roy, D. J. Egger, S. Filipp, F. K. Wilhelm, and S. Machnes, Phys. Rev. Appl. 15, 034080 (2021).
- [25] J. Zhang and M. Sarovar, Phys. Rev. Lett. 113, 080401 (2014).
- [26] A. Sone and P. Cappellaro, Phys. Rev. A 95, 022335 (2017).
- [27] J. Kiukas, K. Yuasa, and D. Burgarth, Phys. Rev. A 95, 052132 (2017).
- [28] D. Burgarth and A. Ajoy, Phys. Rev. Lett. 119, 030402 (2017).
- [29] S.-T. Wang, D.-L. Deng, and L.-M. Duan, New J. Phys. 17, 093017 (2015).
- [30] S. Xue, R. Wu, S. Ma, D. Li, and M. Jiang, Phys. Rev. A 103, 022604 (2021).
- [31] S. Xue, L. Tan, R. Wu, M. Jiang, and I. R. Petersen, Phys. Rev. A 102, 042227 (2020).
- [32] J. Zhang and M. Sarovar, Phys. Rev. A 91, 052121 (2015).
- [33] R. S. Bennink and P. Lougovski, New J. Phys. 21, 083013 (2019).
- [34] L. Ljung, Annu. Rev. Control 34, 1 (2010).
- [35] C. Le Bris, M. Mirrahimi, H. Rabitz, and G. Turinici, ESAIM: COCV 13, 378 (2007).
- [36] O. F. Alis, H. Rabiz, M. Q. Rosenthal, C. Phan, and M. Pence, J. Math. Chem. 35, 65 (2004).
- [37] M. Ndong, J. Salomon, and D. Sugny, J. Phys. A: Math. Theor. 47, 265302 (2014).
- [38] C. R. Rojas, J. S. Welsh, G. C. Goodwin, and A. Feuer, Automatica 43, 993 (2007).
- [39] L. Baudouin and A. Mercado, Appl. Anal. 87, 1145 (2008).
- [40] S. Bonnabel, M. Mirrahimi, and P. Rouchon, Automatica 45, 1144 (2009).
- [41] Y. Fu and G. Turinici, ESAIM: COCV 23, 1129 (2017).
- [42] Y. Wang, D. Dong, B. Qi, J. Zhang, I. R. Petersen, and H. Yonezawa, IEEE Trans. Autom. Control 63, 1388 (2018).
- [43] J.-S. Li and N. Khaneja, Phys. Rev. A 73, 030302(R) (2006).

- [44] J. S. Li and N. Khaneja, IEEE Trans. Auto. Control 54, 528 (2009).
- [45] K. Kobzar, T. E. Skinner, N. Khaneja, S. J. Glaser, and B. Luy, J. Magn. Reson. **194**, 58 (2008).
- [46] M. Lapert, Y. Zhang, M. A. Janich, S. J. Glaser, and D. Sugny, Sci. Rep. 2, 589 (2012).
- [47] G. Turinici, Phys. Rev. A 100, 053403 (2019).
- [48] L. Van Damme, Q. Ansel, S. J. Glaser, and D. Sugny, Phys. Rev. A 95, 063403 (2017).
- [49] A. Ruschhaupt, X. Chen, D. Alonso, and J. G. Muga, New J. Phys. 14, 093040 (2012).
- [50] L. Van-Damme, D. Schraft, G. T. Genov, D. Sugny, T. Halfmann, and S. Guérin, Phys. Rev. A 96, 022309 (2017).
- [51] D. Daems, A. Ruschhaupt, D. Sugny, and S. Guérin, Phys. Rev. Lett. 111, 050404 (2013).
- [52] G. T. Genov, D. Schraft, T. Halfmann, and N. V. Vitanov, Phys. Rev. Lett. **113**, 043001 (2014).
- [53] D. Buterakos, S. Das Sarma, and E. Barnes, PRX Quantum 2, 010341 (2021).
- [54] J. Zeng and E. Barnes, Phys. Rev. A 98, 012301 (2018).
- [55] J. Zeng, C. H. Yang, A. S. Dzurak, and E. Barnes, Phys. Rev. A 99, 052321 (2019).

- [56] Y. Maday and J. Salomon, in *Proceedings of the 48th IEEE Conference on Decision and Control, 2009, Held jointly with the 28th Chinese Control Conference (CDC/CCC 2009)*, Shanghai, China (Institute of Electrical and Electronics Engineers, New York, 2009), pp. 375–379.
- [57] S. Buchwald, G. Ciaramella, and J. Samonon, SIAM J. Control Optim. 59, 4511 (2021).
- [58] S. Conolly, D. Nishimura, and A. Macovski, IEEE Trans. Med. Imag. 5, 106 (1986).
- [59] M. Lapert, Y. Zhang, M. Braun, S. J. Glaser, and D. Sugny, Phys. Rev. Lett. **104**, 083001 (2010).
- [60] M. Levitt, Spin Dynamics: Basics of Nuclear Magnetic Resonance (Wiley, Chichester, UK, 2008).
- [61] T. E. Skinner, T. O. Reiss, B. Luy, N. Khaneja, and S. J. Glaser, J. Magn. Reson. 172, 17 (2005).
- [62] See Supplemental Material at http://link.aps.org/supplemental/ 10.1103/PhysRevA.104.063112 for technical details on the numerical simulations and additional results.
- [63] B. Bonnard, O. Cots, S. J. Glaser, M. Lapert, D. Sugny, and Y. Zhang, IEEE Trans. Autom. Control 57, 1957 (2012).
- [64] Z. Tosner, A. Purea, J. O. Struppe, S. Wegner, F. Engelke, S. J. Glaser, and B. Reif, J. Magn. Reson. 284, 20 (2017)
- [65] F. Engelke, Concepts Magn. Reson. 15, 129 (2002).