

We shall study an underdetermined system of PDEs coming from mathematical physics.

J. Wheeler: Exist potentials for the Einstein equations?

Pommaret (1995): **No.** See J.-F. Pommaret, *Dualité différentielles et applications*, C.R. Acad. Sci. Paris, Serie I 320 (1995). (See also H. Weyl, *Space Time Matter*, fourth edition, Dover, 1952, for more details on Einstein equations.)

Let us prove this by using *OreModules*!

```
> with(Ore_algebra):
> with(OreModules):
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We define the Weyl algebra $\text{Alg} = A_4$ (where $x[1], x[2], x[3]$ stand for the three space components and $x[4] = ct$ for the time):

```
> Alg := DefineOreAlgebra(diff=[D[1],x[1]], diff=[D[2],x[2]],
> diff=[D[3],x[3]], diff=[D[4],x[4]], polynom={x[1],x[2],x[3],x[4]}):
```

The linearized Ricci equations in the vacuum are defined by the following matrix of differential operator.

```
> R := evalm([[D[2]^2+D[3]^2-D[4]^2, D[1]^2, D[1]^2, -D[1]^2, -2*D[1]*D[2], 0, 0,
> -2*D[1]*D[3], 0, 2*D[1]*D[4]],
> [D[2]^2, D[1]^2+D[3]^2-D[4]^2, D[2]^2, -D[2]^2, -2*D[1]*D[2],
> -2*D[2]*D[3], 0, 0, 2*D[2]*D[4], 0],
> [D[3]^2, D[3]^2, D[1]^2+D[2]^2-D[4]^2, -D[3]^2, 0, -2*D[2]*D[3],
> 2*D[3]*D[4], -2*D[1]*D[3], 0, 0],
> [D[4]^2, D[4]^2, D[4]^2, D[1]^2+D[2]^2+D[3]^2, 0, 0, -2*D[3]*D[4], 0,
> -2*D[2]*D[4], -2*D[1]*D[4]],
> [0, 0, D[1]*D[2], -D[1]*D[2], D[3]^2-D[4]^2, -D[1]*D[3], 0,
> -D[2]*D[3], D[1]*D[4], D[2]*D[4]],
> [D[2]*D[3], 0, 0, -D[2]*D[3], -D[1]*D[3], D[1]^2-D[4]^2, D[2]*D[4],
> -D[1]*D[2], D[3]*D[4], 0],
> [D[3]*D[4], D[3]*D[4], 0, 0, 0, -D[2]*D[4], D[1]^2+D[2]^2,
> -D[1]*D[4], -D[2]*D[3], -D[1]*D[3]],
> [0, D[1]*D[3], 0, -D[1]*D[3], -D[2]*D[3], -D[1]*D[2], D[1]*D[4],
> D[2]^2-D[4]^2, 0, D[3]*D[4]],
> [D[2]*D[4], 0, D[2]*D[4], 0, -D[1]*D[4], -D[3]*D[4], -D[2]*D[3], 0,
> D[1]^2+D[3]^2, -D[1]*D[2]],
> [0, D[1]*D[4], D[1]*D[4], 0, -D[2]*D[4], 0, -D[1]*D[3], -D[3]*D[4],
> -D[1]*D[2], D[2]^2+D[3]^2]);
```

$$R := \begin{bmatrix} D_2^2 + D_3^2 - D_4^2, D_1^2, D_1^2, -D_1^2, -2D_1D_2, 0, 0, -2D_1D_3, 0, 2D_1D_4 \\ D_2^2, D_1^2 + D_3^2 - D_4^2, D_2^2, -D_2^2, -2D_1D_2, -2D_2D_3, 0, 0, 2D_2D_4, 0 \\ D_3^2, D_3^2, D_1^2 + D_2^2 - D_4^2, -D_3^2, 0, -2D_2D_3, 2D_3D_4, -2D_1D_3, 0, 0 \\ D_4^2, D_4^2, D_4^2, D_1^2 + D_2^2 + D_3^2, 0, 0, -2D_3D_4, 0, -2D_2D_4, -2D_1D_4 \\ 0, 0, D_1D_2, -D_1D_2, D_3^2 - D_4^2, -D_1D_3, 0, -D_2D_3, D_1D_4, D_2D_4 \\ D_2D_3, 0, 0, -D_2D_3, -D_1D_3, D_1^2 - D_4^2, D_2D_4, -D_1D_2, D_3D_4, 0 \\ D_3D_4, D_3D_4, 0, 0, 0, -D_2D_4, D_1^2 + D_2^2, -D_1D_4, -D_2D_3, -D_1D_3 \\ 0, D_1D_3, 0, -D_1D_3, -D_2D_3, -D_1D_2, D_1D_4, D_2^2 - D_4^2, 0, D_3D_4 \\ D_2D_4, 0, D_2D_4, 0, -D_1D_4, -D_3D_4, -D_2D_3, 0, D_1^2 + D_3^2, -D_1D_2 \\ 0, D_1D_4, D_1D_4, 0, -D_2D_4, 0, -D_1D_3, -D_3D_4, -D_1D_2, D_2^2 + D_3^2 \end{bmatrix}$$

First of all, let us compute the rank of the Alg -module M associated with R .

```
> OreRank(R, Alg);
```

The formal adjoint of the linearized Ricci equations is defined by the following matrix of differential operators.

$$R_{\text{adj}} := \left[\begin{array}{c} D_2^2 + D_3^2 - D_4^2, D_2^2, D_3^2, D_4^2, 0, D_2 D_3, D_3 D_4, 0, D_2 D_4, 0 \\ D_1^2, D_1^2 + D_3^2 - D_4^2, D_3^2, D_4^2, 0, 0, D_3 D_4, D_1 D_3, 0, D_1 D_4 \\ D_1^2, D_2^2, D_1^2 + D_2^2 - D_4^2, D_4^2, D_1 D_2, 0, 0, 0, D_2 D_4, D_1 D_4 \\ -D_1^2, -D_2^2, -D_3^2, D_1^2 + D_2^2 + D_3^2, -D_1 D_2, -D_2 D_3, 0, -D_1 D_3, 0, 0 \\ -2 D_1 D_2, -2 D_1 D_2, 0, 0, D_3^2 - D_4^2, -D_1 D_3, 0, -D_2 D_3, -D_1 D_4, -D_2 D_4 \\ 0, -2 D_2 D_3, -2 D_2 D_3, 0, -D_1 D_3, D_1^2 - D_4^2, -D_2 D_4, -D_1 D_2, -D_3 D_4, 0 \\ 0, 0, 2 D_3 D_4, -2 D_3 D_4, 0, D_2 D_4, D_1^2 + D_2^2, D_1 D_4, -D_2 D_3, -D_1 D_3 \\ -2 D_1 D_3, 0, -2 D_1 D_3, 0, -D_2 D_3, -D_1 D_2, -D_1 D_4, D_2^2 - D_4^2, 0, -D_3 D_4 \\ 0, 2 D_2 D_4, 0, -2 D_2 D_4, D_1 D_4, D_3 D_4, -D_2 D_3, 0, D_1^2 + D_3^2, -D_1 D_2 \\ 2 D_1 D_4, 0, 0, -2 D_1 D_4, D_2 D_4, 0, -D_1 D_3, D_3 D_4, -D_1 D_2, D_2^2 + D_3^2 \end{array} \right]$$

We compute the first extension module ext^1 with values in Alg of the formal adjoint of the linearized Ricci equations in Alg .

$$\begin{aligned} > \text{st} := \text{time}(); \quad \text{Ext1} := \text{Exti}(R_{\text{adj}}, \text{Alg}, 1); \quad \text{time}() - \text{st}; \quad \text{Ext1}[1]; \\ & 44.639 \\ & \left[\begin{array}{c} \%1, 0 \\ 0, \%1, 0 \\ 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0, 0 \\ 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0, 0 \\0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0, 0 \\0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0, 0 \\0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \%1, 0, 0 \\0, \%1, 0, 0 \\0, \%1, 0 \\0, \%1 \end{array} \right] \\ & \%1 := D_1^2 + D_2^2 - D_4^2 + D_3^2 \\ > \text{Ext1}[2]; \end{aligned}$$

$$\left[\begin{array}{l} D_2 D_4, 0, 0, 0, -D_1 D_4, 0, 0, 0, D_1^2, -D_1 D_2 \\ D_3 D_4, 0, 0, 0, 0, 0, D_1^2, -D_1 D_4, 0, -D_1 D_3 \\ D_2 D_3, 0, 0, 0, -D_1 D_3, D_1^2, 0, -D_1 D_2, 0, 0 \\ D_4^2, 0, 0, D_1^2, 0, 0, 0, 0, 0, -2 D_1 D_4 \\ D_3^2, 0, D_1^2, 0, 0, 0, 0, -2 D_1 D_3, 0, 0 \\ D_2^2, D_1^2, 0, 0, -2 D_1 D_2, 0, 0, 0, 0, 0 \\ 0, D_3 D_4, 0, 0, 0, -D_2 D_4, D_2^2, 0, -D_2 D_3, 0 \\ 0, D_4^2, 0, D_2^2, 0, 0, 0, 0, -2 D_2 D_4, 0 \\ 0, D_3^2, D_2^2, 0, 0, -2 D_2 D_3, 0, 0, 0, 0 \\ 0, -D_1 D_4, 0, 0, D_2 D_4, 0, 0, 0, D_1 D_2, -D_2^2 \\ 0, -D_1 D_3, 0, 0, D_2 D_3, D_1 D_2, 0, -D_2^2, 0, 0 \\ 0, 0, D_4^2, D_3^2, 0, 0, -2 D_3 D_4, 0, 0, 0 \\ 0, 0, -D_2 D_4, 0, 0, D_3 D_4, D_2 D_3, 0, -D_3^2, 0 \\ 0, 0, -D_1 D_4, 0, 0, 0, D_1 D_3, D_3 D_4, 0, -D_3^2 \\ 0, 0, D_1 D_2, 0, D_3^2, -D_1 D_3, 0, -D_2 D_3, 0, 0 \\ 0, 0, 0, D_2 D_3, 0, D_4^2, -D_2 D_4, 0, -D_3 D_4, 0 \\ 0, 0, 0, D_1 D_3, 0, 0, -D_1 D_4, D_4^2, 0, -D_3 D_4 \\ 0, 0, 0, D_1 D_2, D_4^2, 0, 0, 0, -D_1 D_4, -D_2 D_4 \\ 0, 0, 0, 0, D_3 D_4, -D_1 D_4, D_1 D_2, 0, 0, -D_2 D_3 \\ 0, 0, 0, 0, 0, -D_1 D_4, 0, D_2 D_4, D_1 D_3, -D_2 D_3 \end{array} \right]$$

Interpretation of the result:

The second matrix $\text{Ext1}[2]$ gives a family of generators of 20 torsion elements in the Alg -module M associated with the linearized Ricci equations. The first matrix $\text{Ext1}[1]$ shows that all the torsion elements are killed by the same differential operator, namely the Dalembertian $D_1^2 - D_4^2 + D_2^2 + D_3^2$. The third matrix gives a parametrization of the torsion-free module $M / t(M)$.

In conclusion, the linearized Ricci equations cannot be parametrized, i.e., they do not admit a potential.

Let us rewrite the torsion elements in terms of the system variable y_1, \dots, y_{10} :

```
> TorsionElements(R, [seq(y[i](x[1],x[2],x[3],x[4]), i=1..10)], Alg);
```

$$\%1 := \theta_{20}(x_1, x_2, x_3, x_4)$$

$$\%2 := \theta_{19}(x_1, x_2, x_3, x_4)$$

$$\%3 := \theta_{18}(x_1, x_2, x_3, x_4)$$

$$\%4 := \theta_{17}(x_1, x_2, x_3, x_4)$$

$$\%5 := \theta_{16}(x_1, x_2, x_3, x_4)$$

$$\%6 := \theta_{15}(x_1, x_2, x_3, x_4)$$

$$\%7 := \theta_{14}(x_1, x_2, x_3, x_4)$$

%8 := $\theta_{13}(x_1, x_2, x_3, x_4)$

%9 := $\theta_{12}(x_1, x_2, x_3, x_4)$

%10 := $\theta_{11}(x_1, x_2, x_3, x_4)$

%11 := $\theta_{10}(x_1, x_2, x_3, x_4)$

$\%12 := \theta_9(x_1, x_2, x_3, x_4)$

%13 := $\theta_8(x_1, x_2, x_3, x_4)$

%14 := $\theta_7(x_1, x_2, x_3, x_4)$

%15 := $\theta_6(x_1 \ x_2 \ x_3 \ x_4)$

%16 := $\theta_5(x_1, x_2, x_3, x_4)$

%17 := $\theta_4(x_1, x_2, x_3, x_4)$

$$\%18 := \theta_3(x_1, x_2, x_3, x_4)$$

%19 := $\theta_3(x_1, x_2, x_3, x_4)$

%20 := $\theta_1(x_1, x_2, x_3, x_4)$

$$\omega_1(\omega_1, \omega_2, \omega_3, \omega_4)$$

$$\begin{aligned}
\theta_1(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_4 \partial x_2} \%10) - (\frac{\partial^2}{\partial x_4 \partial x_1} \%6) + (\frac{\partial^2}{\partial x_1^2} \%2) - (\frac{\partial^2}{\partial x_2 \partial x_1} \%1) \\
\theta_2(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_4 \partial x_3} \%10) + (\frac{\partial^2}{\partial x_1^2} \%5) - (\frac{\partial^2}{\partial x_4 \partial x_1} \%3) - (\frac{\partial^2}{\partial x_3 \partial x_1} \%1) \\
\theta_3(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_3 \partial x_2} \%10) - (\frac{\partial^2}{\partial x_3 \partial x_1} \%6) + (\frac{\partial^2}{\partial x_1^2} \%4) - (\frac{\partial^2}{\partial x_2 \partial x_1} \%3) \\
\theta_4(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_1^2} \%10) + (\frac{\partial^2}{\partial x_1^2} \%7) - 2(\frac{\partial^2}{\partial x_4 \partial x_1} \%1) \\
\theta_5(x_1, x_2, x_3, x_4) &= (\frac{\partial^3}{\partial x_2^2} \%10) + (\frac{\partial^2}{\partial x_1^2} \%8) - 2(\frac{\partial^2}{\partial x_3 \partial x_1} \%3) \\
\theta_6(x_1, x_2, x_3, x_4) &= (\frac{\partial^3}{\partial x_2^2} \%10) + (\frac{\partial^2}{\partial x_1^2} \%9) - 2(\frac{\partial^2}{\partial x_2 \partial x_1} \%6) \\
\theta_7(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_4 \partial x_3} \%9) - (\frac{\partial^2}{\partial x_4 \partial x_2} \%4) + (\frac{\partial^2}{\partial x_2^2} \%5) - (\frac{\partial^2}{\partial x_3 \partial x_2} \%2) \\
\theta_8(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_1^2} \%9) + (\frac{\partial^2}{\partial x_2^2} \%7) - 2(\frac{\partial^2}{\partial x_4 \partial x_2} \%2) \\
\theta_9(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_3^2} \%9) + (\frac{\partial^2}{\partial x_2^2} \%8) - 2(\frac{\partial^2}{\partial x_3 \partial x_2} \%4) \\
\theta_{10}(x_1, x_2, x_3, x_4) &= -(\frac{\partial^2}{\partial x_4 \partial x_1} \%9) + (\frac{\partial^2}{\partial x_4 \partial x_2} \%6) + (\frac{\partial^2}{\partial x_2 \partial x_1} \%2) - (\frac{\partial^2}{\partial x_2^2} \%1) \\
\theta_{11}(x_1, x_2, x_3, x_4) &= -(\frac{\partial^2}{\partial x_3 \partial x_1} \%9) + (\frac{\partial^2}{\partial x_3 \partial x_2} \%6) + (\frac{\partial^2}{\partial x_2 \partial x_1} \%4) - (\frac{\partial^2}{\partial x_2^2} \%3) \\
\theta_{12}(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_4^2} \%8) + (\frac{\partial^2}{\partial x_3^2} \%7) - 2(\frac{\partial^2}{\partial x_4 \partial x_3} \%5) \\
\theta_{13}(x_1, x_2, x_3, x_4) &= -(\frac{\partial^2}{\partial x_4 \partial x_2} \%8) + (\frac{\partial^2}{\partial x_4 \partial x_3} \%4) + (\frac{\partial^2}{\partial x_3 \partial x_2} \%5) - (\frac{\partial^2}{\partial x_3^2} \%2) \\
\theta_{14}(x_1, x_2, x_3, x_4) &= -(\frac{\partial^2}{\partial x_4 \partial x_1} \%8) + (\frac{\partial^2}{\partial x_3 \partial x_1} \%5) + (\frac{\partial^2}{\partial x_4 \partial x_3} \%3) - (\frac{\partial^2}{\partial x_3^2} \%1) \\
\theta_{15}(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_2^2} \%8) + (\frac{\partial^2}{\partial x_3^2} \%6) - (\frac{\partial^2}{\partial x_3 \partial x_1} \%4) - (\frac{\partial^2}{\partial x_3 \partial x_2} \%3) \\
\theta_{16}(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_3 \partial x_2} \%7) + (\frac{\partial^2}{\partial x_4^2} \%4) - (\frac{\partial^2}{\partial x_4 \partial x_2} \%5) - (\frac{\partial^2}{\partial x_4 \partial x_3} \%2) \\
\theta_{17}(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_3 \partial x_1} \%7) - (\frac{\partial^2}{\partial x_4 \partial x_1} \%5) + (\frac{\partial^2}{\partial x_4^2} \%3) - (\frac{\partial^2}{\partial x_4 \partial x_3} \%1) \\
\theta_{18}(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_2^2} \%7) + (\frac{\partial^2}{\partial x_4^2} \%6) - (\frac{\partial^2}{\partial x_4 \partial x_1} \%2) - (\frac{\partial^2}{\partial x_4 \partial x_2} \%1) \\
\theta_{19}(x_1, x_2, x_3, x_4) &= (\frac{\partial^2}{\partial x_4 \partial x_3} \%6) - (\frac{\partial^2}{\partial x_4 \partial x_1} \%4) + (\frac{\partial^2}{\partial x_2 \partial x_1} \%5) - (\frac{\partial^2}{\partial x_3 \partial x_2} \%1) \\
\theta_{20}(x_1, x_2, x_3, x_4) &= -(\frac{\partial^2}{\partial x_4 \partial x_1} \%4) + (\frac{\partial^2}{\partial x_4 \partial x_2} \%3) + (\frac{\partial^2}{\partial x_3 \partial x_1} \%2) - (\frac{\partial^2}{\partial x_3 \partial x_2} \%1)
\end{aligned}$$

```

%1 := y10(x1, x2, x3, x4)
%2 := y9(x1, x2, x3, x4)
%3 := y8(x1, x2, x3, x4)
%4 := y6(x1, x2, x3, x4)
%5 := y7(x1, x2, x3, x4)
%6 := y5(x1, x2, x3, x4)
%7 := y4(x1, x2, x3, x4)
%8 := y3(x1, x2, x3, x4)
%9 := y2(x1, x2, x3, x4)
%10 := y1(x1, x2, x3, x4)

```

The parametrization of the torsion-free *Alg*-module $M / \text{t}(M)$ is defined by:

```
> Ext1[3];
```

$$\begin{bmatrix} -2D_1 & 0 & 0 & 0 \\ 0 & -2D_2 & 0 & 0 \\ 0 & 0 & -2D_3 & 0 \\ 0 & 0 & 0 & -2D_4 \\ -D_2 & -D_1 & 0 & 0 \\ 0 & -D_3 & -D_2 & 0 \\ 0 & 0 & -D_4 & -D_3 \\ -D_3 & 0 & -D_1 & 0 \\ 0 & -D_4 & 0 & -D_2 \\ -D_4 & 0 & 0 & -D_1 \end{bmatrix}$$

Let us point out that this parametrization $Ext1[3]$ depends on 4 arbitrary potentials. As the rank of the Alg -module M is also 4, and thus, the rank of the Alg -module $M / t(M)$ is also 4, we deduce that $Ext1[3]$ is a minimal parametrization of $M / t(M)$.

We compute the second extension module ext^2 with values in Alg of the formal adjoint of the linearized Ricci equations in Alg . The interpretation is the same as before.

```
> st := time(): Ext2 := Exti(R_adj, Alg, 2): time() - st; Ext2[1];
2.839

$$\begin{bmatrix} D_1 & 0 & 0 & 0 \\ D_4^2 & 0 & 0 & 0 \\ D_3 D_4 & 0 & 0 & 0 \\ D_2 D_4 & 0 & 0 & 0 \\ D_3^2 & 0 & 0 & 0 \\ D_2 D_3 & 0 & 0 & 0 \\ D_2^2 & 0 & 0 & 0 \\ 0 & D_2 & 0 & 0 \\ 0 & D_4^2 & 0 & 0 \\ 0 & D_3 D_4 & 0 & 0 \\ 0 & D_1 D_4 & 0 & 0 \\ 0 & D_3^2 & 0 & 0 \\ 0 & D_1 D_3 & 0 & 0 \\ 0 & D_1^2 & 0 & 0 \\ 0 & 0 & D_3 & 0 \\ 0 & 0 & D_4^2 & 0 \\ 0 & 0 & D_2 D_4 & 0 \\ 0 & 0 & D_1 D_4 & 0 \\ 0 & 0 & D_2^2 & 0 \\ 0 & 0 & D_1 D_2 & 0 \\ 0 & 0 & D_1^2 & 0 \\ 0 & 0 & 0 & D_4 \\ 0 & 0 & 0 & D_3^2 \\ 0 & 0 & 0 & D_2 D_3 \\ 0 & 0 & 0 & D_1 D_3 \\ 0 & 0 & 0 & D_2^2 \\ 0 & 0 & 0 & D_1 D_2 \\ 0 & 0 & 0 & D_1^2 \end{bmatrix}$$

```

We compute a free resolution of the linearized Ricci equations.

```
> st:=time(): FreeResolution(R, Alg); time()-st;
```

```

table([1 =

$$\begin{bmatrix} D_2^2 + D_3^2 - D_4^2, D_1^2, D_1^2, -D_1^2, -2D_1 D_2, 0, 0, -2D_1 D_3, 0, 2D_1 D_4 \\ D_2^2, D_1^2 + D_3^2 - D_4^2, D_2^2, -D_2^2, -2D_1 D_2, -2D_2 D_3, 0, 0, 2D_2 D_4, 0 \\ D_3^2, D_3^2, D_1^2 + D_2^2 - D_4^2, -D_3^2, 0, -2D_2 D_3, 2D_3 D_4, -2D_1 D_3, 0, 0 \\ D_4^2, D_4^2, D_4^2, D_1^2 + D_2^2 + D_3^2, 0, 0, -2D_3 D_4, 0, -2D_2 D_4, -2D_1 D_4 \\ 0, 0, D_1 D_2, -D_1 D_2, D_3^2 - D_4^2, -D_1 D_3, 0, -D_2 D_3, D_1 D_4, D_2 D_4 \\ D_2 D_3, 0, 0, -D_2 D_3, -D_1 D_3, D_1^2 - D_4^2, D_2 D_4, -D_1 D_2, D_3 D_4, 0 \\ D_3 D_4, D_3 D_4, 0, 0, 0, -D_2 D_4, D_1^2 + D_2^2, -D_1 D_4, -D_2 D_3, -D_1 D_3 \\ 0, D_1 D_3, 0, -D_1 D_3, -D_2 D_3, -D_1 D_2, D_1 D_4, D_2^2 - D_4^2, 0, D_3 D_4 \\ D_2 D_4, 0, D_2 D_4, 0, -D_1 D_4, -D_3 D_4, -D_2 D_3, 0, D_1^2 + D_3^2, -D_1 D_2 \\ 0, D_1 D_4, D_1 D_4, 0, -D_2 D_4, 0, -D_1 D_3, -D_3 D_4, -D_1 D_2, D_2^2 + D_3^2 \end{bmatrix},$$

2 = 
$$\begin{bmatrix} -D_4 & -D_4 & -D_4 & -D_4 & 0 & 0 & 2D_3 & 0 & 2D_2 & 2D_1 \\ -D_3 & -D_3 & D_3 & D_3 & 0 & 2D_2 & -2D_4 & 2D_1 & 0 & 0 \\ -D_2 & D_2 & -D_2 & D_2 & 2D_1 & 2D_3 & 0 & 0 & -2D_4 & 0 \\ D_1 & -D_1 & -D_1 & D_1 & 2D_2 & 0 & 0 & 2D_3 & 0 & -2D_4 \end{bmatrix},$$

3 = INJ(4)
])
4.190

```

We compute a free resolution of the formal adjoint of the linearized Ricci equations.

```

> st:=time(): FreeResolution(R_adj, Alg); time()-st;





```